

## Spectrum of hot water in the 2000–4750 $\text{cm}^{-1}$ frequency range

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### Abstract

An emission spectrum recorded in an oxyacetylene torch [P.-F. Coheur, P.F. Bernath, M. Carleer, R. Colin, O.L. Polyansky, N.F. Zobov, S.V. Shirin, R.J. Barber, J. Tennyson, *J. Chem. Phys.* 122 (2005) 074307] is analyzed for the region covering stretching fundamentals and associated hot bands of water. Many lines could be assigned on the basis of previously determined energy levels. New assignments made with a new variational linelist allow a further 800 energy levels covering 15 vibrational states and rotations up to  $J = 32$  to be assigned. A simultaneous re-analysis of previously reported sunspot absorption spectra leads to the assignment of 581 further lines in the  $L$ -band spectrum and 67 in the  $N$ -band spectrum.

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### 1. Introduction

The spectroscopy of water vapor finds application in a wide variety of areas including combustion science, atmospheric remote sensing, and astronomy [1]. In particular, hot water vapor is a prominent contributor to the spectral energy distributions of late M dwarf stars [2] and large sunspots [3,4]. These astronomical objects have effective temperatures of about 3000 K. The acquisition and interpretation of laboratory spectra of water vapor at 3000 K are therefore required for the calculation of molecular opacities [5] that are used to simulate these astronomical spectra. In fact the required molecular opacities are

probably best calculated from some combination of energy levels and line intensities obtained from experiment with the predictions of state-of-the-art quantum chemistry [6].

The first modern high resolution spectra of hot water vapor were recorded by Maillard at the Meudon Observatory in France. He recorded emission spectra of an oxyacetylene torch with a Fourier transform spectrometer and the analysis of these spectra was an important landmark in the spectroscopy of water vapor [7,8]. The advantage of the oxy-acetylene torch is that it produces a spectrum of water vapor at 3000 K, although the lines are nearly  $0.1 \text{ cm}^{-1}$  wide primarily due to pressure-broadening at ambient atmospheric pressure. We have repeated the Maillard torch emission experiment and obtained new data over a wider spectral range and with somewhat higher signal-to-noise ratio. In our first paper based on these new data [9], we reported on the analysis of the pure rotation and  $\nu_2$  bending mode region, 500–2000  $\text{cm}^{-1}$ . The bending mode region yielded energy levels up to  $9\nu_2$ , which is over

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the barrier to linearity and provided evidence for “monodromy” [10]. In the work reported below, we extend our analysis to the  $\nu_1$  and  $\nu_3$  stretching mode region, 2000–4750  $\text{cm}^{-1}$ .

## 2. Observed spectra

The laboratory spectra reported by Coheur et al. [9] were obtained by recording emission from an oxy-acetylene torch using a Bruker IFS 120 M Fourier transform spectrometer between 500 and 13000  $\text{cm}^{-1}$ . In the region considered here (2000–4750  $\text{cm}^{-1}$ ) an InSb detector was used with a  $\text{CaF}_2$  window and beamsplitter. Either a 2 or a 4 mm aperture was chosen and the spectral resolution was set to 0.05  $\text{cm}^{-1}$  (18 cm maximum optical path difference). Five hundred and twelve scans were co-added, thereby producing emission spectra with very low noise. Spectra were recorded with the torch at atmospheric pressure which, combined with a temperature of about 3000 K, leads to broadened lines with many blends and an uncertainty of about 0.02  $\text{cm}^{-1}$  in the determination of line positions. A few lines have fitting errors larger than this value; information on this is given in the archived version of the spectrum.

As is obvious from Fig. 1, the spectra are very dense, showing, in addition to water lines, emission features of CO,  $\text{CO}_2$  and OH. The OH lines were identified in the spectra by using the spectroscopic constants of Colin et al. [11]. Vibrational levels up to  $v = 8$  were considered in the comparison, and very weak satellite lines were neglected. Frequencies of OH rotational transitions were taken from Mélen et al. [12]. CO lines were similarly identified by using the HITEMP database [13] as a reference, with  $v = 8$  as the highest vibrational level, and both the  $^{12}\text{CO}$  and  $^{13}\text{CO}$

isotopes were considered. Many of the possible OH and CO lines identified in the spectrum and marked in the line list are likely due to water. These transitions match CO or OH line positions but only for transitions with higher vibrational or rotational quantum numbers than expected for a 3000 K source.

Carbon dioxide transitions form a major part of the spectrum: the region between 2200 and 2400  $\text{cm}^{-1}$  is dominated by  $\text{CO}_2$  lines and was not analyzed. Outside this region  $\text{CO}_2$  lines were identified using data taken from the carbon dioxide spectroscopic database (CDS) system of the Institute of Atmospheric Optics of the Siberian Branch the Russian Academy of Sciences (see <http://spectra.iao.ru>).

15229 lines were measured in the 2000–2200 and 2400–4750  $\text{cm}^{-1}$  regions, although, as discussed below, blends mean that these correspond to 19008 transitions.

Two absorption spectra of sunspots [14] were reconsidered at the same time as the laboratory emission spectra. These sunspot spectra have an estimated water vapor temperature of 3200 K. The first covered the *L*-band (2497–3195  $\text{cm}^{-1}$ ), which had been previously analyzed [15], and contains 2723 lines of which 1207 have been assigned [15]. The second was the *N*-band spectrum spanning the 722–1011  $\text{cm}^{-1}$  region [3], which has also been partially analyzed before [4,9,10,16]. Obviously the *N*-band spectrum does not overlap the spectral regions considered here, but it is extremely rich and many of the new assignments made in the laboratory spectrum could be confirmed by transitions in *N*-band spectrum using predicted combination differences.

Fig. 1 compares a small portion of the laboratory emission spectrum with the sunspot absorption spectrum of Wallace et al. [14]. Table 1 presents the corresponding

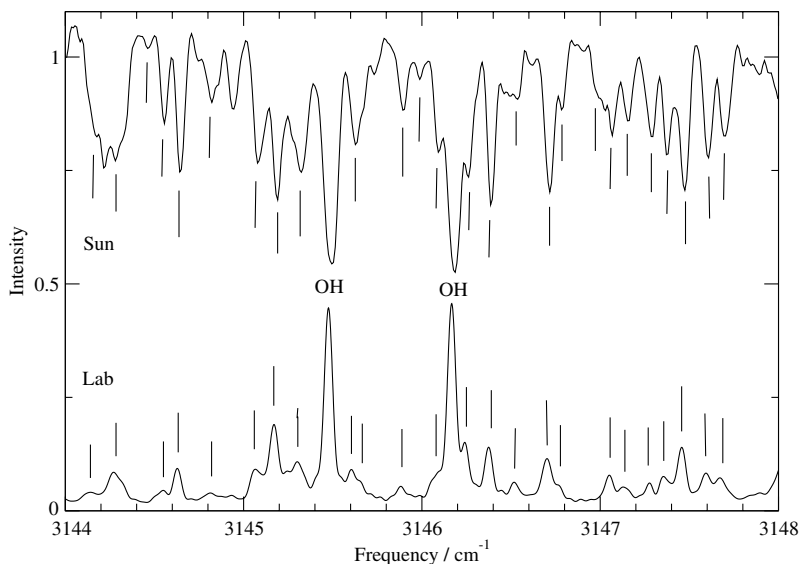


Fig. 1. Comparison of a sunspot absorption spectrum (upper trace) [14] and the laboratory emission spectrum (lower trace). The vertical lines indicate the location of assigned transitions as reported in Table 1. The vertical scale on the left refers to the sunspot absorption spectrum; the laboratory emission spectrum was not obtained on an absolute scale.

Table 1  
Laboratory emission linelist in the region 3144–3148 cm<sup>-1</sup>, corresponding to Fig. 1

Wavenumber (cm <sup>-1</sup> )	Peak height arb. units	$J'$	$K'_a$	$K'_c$	$J''$	$K''_a$	$K''_c$	$v'_1v'_2v'_3 - v''_1v''_2v''_3$	Notes
3144.15772	0.0166								OH
3144.29116	0.0557	14	10	4	15	10	5	101 – 100	dnc
3144.55555	0.0217	9	0	9	10	0	10	300 – 101	d
		9	1	8	10	3	7	022 – 021	
3144.64157	0.0721	22	3	20	23	3	21	002 – 001	nc
		19	7	12	20	7	13	002 – 001	
3144.82506	0.0104	8	2	7	9	2	8	300 – 101	
		9	3	7	10	3	8	211 – 210	
3144.93731	0.0132								
3145.08231	0.0456	17	4	14	18	4	15	101 – 100	
3145.18235	0.1500	22	5	17	23	5	18	001 – 000	
3145.31097	0.0523	22	2	21	23	2	22	012 – 011	nc
3145.31097	0.0523	11	3	8	12	5	7	012 – 011	
3145.48902	0.4270								OH
3145.61306	0.0345	15	6	9	16	5	12	030 – 010	
		12	2	11	13	2	12	200 – 001	
3145.65105	0.0236	16	9	7	17	9	8	021 – 100	n
		17	3	14	18	3	15	101 – 100	
3145.89319	0.0192	11	3	9	12	5	8	001 – 000	
3146.08289	0.0299	21	5	16	22	5	17	011 – 010	
3146.18003	0.4620								OH
3146.25914	0.0779	24	1	24	25	1	25	002 – 001	d
3146.38747	0.1040	20	0	20	21	0	21	101 – 100	d
		11	2	10	12	2	11	201 – 200	
3146.53177	0.0255	20	2	19	21	1	20	110 – 010	
		14	5	9	15	6	10	100 – 000	
3146.71496	0.0861	16	5	11	17	5	12	101 – 100	
3146.78179	0.0242	20	7	14	21	7	15	011 – 010	
		18	4	14	19	5	15	100 – 000	
3147.06398	0.0482	19	3	16	20	4	17	100 – 000	nc
		24	1	24	25	1	25	031 – 030	
3147.14849	0.0214	21	4	18	22	4	19	002 – 001	
3147.28671	0.0332	24	0	24	25	0	25	031 – 030	
3147.37491	0.0400	15	5	11	16	5	12	111 – 110	
3147.46963	0.1040	21	3	18	22	3	19	002 – 001	d
		25	1	25	26	1	26	021 – 020	
3147.60458	0.0494	20	5	15	21	5	16	021 – 020	
		19	2	17	20	3	18	110 – 010	
3147.68407	0.0345	11	10	1	12	11	2	002 – 100	dnc

The notes mean: ‘n,’ new upper level; ‘c,’ confirmed by combination differences; ‘d,’ doublet due to degeneracy. Note the double assignment of many of the lines.

emission spectrum linelist for this small portion of the spectrum only. The full linelist as well as new versions of the sunspot spectra with updated assignments can be found in the electronic archive.

### 3. Line assignments

Line assignments were performed using the BT2 linelist [5]. This linelist was computed in order to model the spectrum of water vapor in cool stars and is thus suitable for the temperatures considered here. BT2 contains all transitions between levels of water with  $J \leq 50$  and energy less than 30 000 cm<sup>-1</sup> above the  $J = 0$  ground state. Transition wavenumbers and wavefunctions were calculated using the spectroscopically determined potential energy surface of Shirin et al. [18]; transition intensities were obtained using the dipole moment surface of Schwenke and Partridge [19].

Table 2  
Summary of new energy levels

Vib. level	Number of new levels	Highest $J$	
		Prev.	This work
(100)	62	30	36
(001)	65	32	37
(110)	80	27	32
(011)	28	30	35
(120)	31	20	23
(021)	25	30	32
(200)	74	19	29
(101)	37	28	29
(002)	137	26	32
(130)	15	14	14
(031)	36	25	29
(111)	16	26	22
(012)	55	20	29
(041)	15	25	25
(121)	20	15	20
(022)	19	19	26
(131)	14	11	22

Table 3  
Newly assigned transitions in the sunspot *N*-band spectrum

Wavenumber	$J'$	$K'_a$	$K'_c$	$J''$	$K''_a$	$K''_c$	$v'_1 v'_2 v'_3 - v''_1 v''_2 v''_3$
733.63047	20	7	13	19	6	14	021 – 021
738.78441	20	10	11	19	9	10	110 – 030
747.22399	20	7	13	19	6	14	031 – 031
760.93078	21	11	10	20	10	11	110 – 110
764.85815	19	6	14	18	6	13	021 – 040
767.36292	20	16	4	19	15	5	011 – 011
768.78462	14	6	9	13	3	10	120 – 120
774.16285	18	7	12	17	4	13	002 – 002
777.89056	24	10	15	23	9	14	100 – 020
778.32496	20	20	1	19	19	0	110 – 110
793.14810	20	5	16	19	10	9	120 – 030
797.86795	17	6	12	16	3	13	110 – 110
805.93601	16	8	9	15	7	8	200 – 040
809.74188	14	4	11	13	1	12	130 – 130
819.68444	22	16	7	21	15	6	110 – 110
820.86278	16	5	12	15	2	13	120 – 120
825.22767	22	18	5	21	17	4	110 – 110
826.31982	17	4	13	16	3	14	120 – 120
826.86678	17	3	14	16	2	15	002 – 002
827.50499	23	14	9	22	13	10	110 – 110
836.59448	15	3	13	14	0	14	110 – 110
837.15424	18	5	14	17	2	15	200 – 200
840.18222	23	16	7	22	15	8	110 – 110
840.94358	16	3	13	15	2	14	110 – 110
843.02101	24	17	7	23	16	8	001 – 001
843.96547	23	17	6	22	16	7	110 – 110
853.42384	20	8	13	19	7	12	110 – 030
854.06118	24	15	10	23	14	9	110 – 110
867.38053	25	18	8	24	17	7	001 – 001
871.75264	19	5	14	18	4	15	021 – 021
872.10943	22	8	15	21	7	14	110 – 030
875.49814	21	6	15	20	5	16	110 – 110
880.67718	19	4	15	18	3	16	002 – 002
883.57597	21	6	15	20	5	16	001 – 001
890.45791	26	19	7	25	18	8	001 – 001
896.91056	23	7	16	22	6	17	100 – 100
898.79726	23	8	16	24	9	15	020 – 010
901.01798	30	0	30	30	1	29	021 – 011
914.84521	29	1	29	29	2	28	021 – 011
916.47109	29	11	18	28	10	19	010 – 010
923.31000	19	6	14	18	3	15	021 – 021
926.75156	22	6	16	21	5	17	100 – 100
927.30019	21	5	16	20	4	17	002 – 002
931.85413	29	10	19	28	9	20	010 – 010
933.06676	22	6	16	21	5	17	020 – 020
934.69998	18	8	11	17	5	12	110 – 110
947.65682	19	6	14	20	8	13	021 – 030
950.33241	17	3	15	16	0	16	110 – 110
952.50148	17	3	14	16	2	15	120 – 120
954.56809	16	3	14	15	0	15	120 – 120
957.91102	22	6	16	21	5	17	021 – 021
960.31126	20	8	13	19	5	14	110 – 110
961.85750	27	19	8	26	18	9	020 – 020
962.20082	21	5	16	20	4	17	110 – 110
963.78980	31	1	31	31	2	30	011 – 001
965.80308	23	6	17	22	5	18	002 – 002
966.77900	23	8	16	22	5	17	100 – 100
968.56967	30	14	17	29	13	16	010 – 010
974.10225	16	5	12	17	6	11	120 – 110
975.26034	22	6	17	21	3	18	002 – 002
976.04084	22	10	13	23	11	12	110 – 100
976.74869	20	4	17	19	1	18	200 – 200
981.15597	30	1	30	30	2	29	110 – 100
986.31764	10	2	9	11	3	8	130 – 120

Table 3 (continued)

Wavenumber	$J'$	$K'_a$	$K'_c$	$J''$	$K''_a$	$K''_c$	$v'_1v'_2v'_3 - v''_1v''_2v''_3$
988.17414	18	4	15	17	1	16	120 – 120
1008.38530	22	8	15	21	5	16	110 – 110
1010.06112	20	9	12	21	10	11	110 – 020

The first step in analyzing the torch spectrum was to mark transitions belonging to other species: we firmly identify 389 OH lines and 825 CO lines; a further 984 lines are identified as possibly belonging to OH or CO transitions. 1180 CO<sub>2</sub> lines are identified in 3400–3800 cm<sup>-1</sup> region.

The second step was to make trivial assignments—that is assignments that can be made using previously determined experimental energy levels [20]. Simple checks were made using the theoretical intensities to ensure that the intensity of each assigned transition was consistent with observation. It was possible to assign 12431 lines in this fashion. At this stage 2965 lines remained unassigned.

Two different methods were used to assign new lines. For transitions involving states with high  $J$  but low  $K_a$ , the method of branches [16,21] was used to follow a series of transitions with quantum numbers which simply differ by one in  $J$ . This method was used for 18 low-lying vibrational levels and yielded about 300 new energy levels.

For transitions with  $J$  in the range of 15–25 and intermediate values of  $K_a$ , the error (obs. – calc.) obtained using the BT2 predictions varies smoothly with  $K_a$  and  $K_c$ . This means that the positions of experimentally determined energy levels could be predicted with sufficient accuracy for new assignments to be made. This method was used for 15 vibrational states associated with stretching transitions ( $\Delta v_1$  or  $\Delta v_3 = 1$ ) to yield about 500 new energy levels.

A copy of the assigned lines from the laboratory spectrum for the 2000–4750 cm<sup>-1</sup> region has been placed in the electronic archive. Table 2 summarizes the information obtained for the main vibrational states analyzed; there are a number of other vibrational states for which a few (less than 10) new levels were obtained. It can be seen that our analysis has extended the range of rotational excitation for a significant number of vibrational states. The greatest number of new energy levels belong to the (002) state; transitions involving this level appear to be under-represented in the analysis of previous hot water spectra.

Most of the transitions in the 2000–4750 cm<sup>-1</sup> region involve changes in one of the stretching quantum numbers, either  $v_1$  or  $v_3$ , by one quantum. However, at higher wavenumbers there are also transitions which correspond to the P-branch of two quantum transitions involving a simultaneous change in a stretching quantum number and the bending quantum number  $v_2$ . It is to be expected that further transitions of this nature will be identified at higher frequencies.

Copies of the sunspot absorption spectra with our new assignments have also been placed in the electronic archive. The  $L$ -band spectrum contains only 2723 transitions as against 4593 in the torch emission spectrum in the same region and contains only a few new lines. New assignments

to the sunspot spectrum were therefore largely taken from those made for the torch spectrum. This process resulted in the assignment of 1579 water transitions, 199 OH lines and 10 HCl lines; a total of 1788 assignments as opposed to the 1207 made previously [15].

The  $N$ -band sunspot spectrum is very dense [4,16] and the present study just focused on a few transitions which could be used to confirm assignments in the torch spectrum. As a result of this 67 new transitions were assigned; these are listed in Table 3.

Table 4 presents the new energy levels obtained from our analysis of the spectra discussed above. In this table, we concentrate on high rotational quantum numbers for which there is no previous information. We estimate the error for each of these levels to be about 0.02 cm<sup>-1</sup>, similar to the error in determining the line positions in the laboratory emission spectrum. The table indicates the levels which have been confirmed by combination differences. The values for the other levels, which all rely on a single assigned transition, must be regarded as less secure. However, our experience has shown that the methods we employ are reliable and it is likely that the vast majority of these unconfirmed levels are also correct.

#### 4. Discussion and conclusion

By the standards of modern Fourier transform spectroscopy, the oxy-acetylene torch produces spectra of only modest resolution because of the extensive pressure-broadening of the lines. In addition, Doppler-broadening also makes a contribution (0.09 cm<sup>-1</sup> at 10000 cm<sup>-1</sup> for 3000 K) to the typical linewidth of about 0.1 cm<sup>-1</sup> above 5000 cm<sup>-1</sup>. Below 5000 cm<sup>-1</sup> our average linewidth is about 0.06 cm<sup>-1</sup> and nearly all of the lines are overlapped, see Table 1. The new BT2 linelist [5], however, predicts energy levels to between 0.2 and 0.4 cm<sup>-1</sup> for  $J = 30$ . However, there is considerable cancellation of errors for transition wavenumbers for which this error is more typically 0.1 cm<sup>-1</sup>, a figure which can be very significantly reduced by employing the method of branches discussed above. Furthermore the BT2 linelist has reliable line intensities.

Although our new spectra cover a wider spectral range and have a better signal-to-noise ratio than the earlier Meudon spectra, it is the availability of the BT2 linelist that makes assignments possible. In the stretching mode region (2000–4750 cm<sup>-1</sup>) covered in this paper, the torch spectra have provided extensive new highly excited vibration-rotation energy levels for more than a dozen vibrations (Table 2). In addition to the stretching fundamental (100) and (001), major improvements have been made to the vibrational levels that have two quanta of stretch, (002) and

Table 4  
New energy levels, in  $\text{cm}^{-1}$ , for high  $J$  vibration–rotation levels

$J$	$K_a$	$K_c$	Wavenumber
(100)			
31	0	31	12748.5829
31	1	31	12748.5829
31	1	30	13291.7717
31	2	30	13291.7717
31	2	29	13759.1605
31	3	29	13759.1605
31	3	28	14198.9484
31	4	28	14198.9484
31	4	27	14601.7860
31	19	13	18848.7981
31	19	12	18848.7981
32	0	32	13310.1642
32	1	32	13310.1642
32	2	31	13870.8937
32	2	30	14346.0170
32	3	30	14346.0170
32	20	13	19848.2666
32	20	12	19848.2666
33	0	33	13887.1749
33	1	33	13887.1749
33	1	32	14462.3518
34	1	33	15070.5999
34	2	33	15070.5999
35	0	35	15086.5589
35	1	35	15086.5589
35	1	34	15694.1713
35	2	34	15694.1713
36	0	36	15708.5886
36	1	36	15708.5886
(001)			
32	2	30	14474.2885 <sup>a</sup>
32	3	30	14474.2885 <sup>a</sup>
32	3	29	14921.4139 <sup>a</sup>
32	4	29	14921.4139 <sup>a</sup>
32	5	27	15713.6124
32	6	27	15713.6124
33	0	33	14023.1636
33	1	33	14023.1636
33	1	32	14593.1539 <sup>a</sup>
33	2	32	14593.1539 <sup>a</sup>
33	2	31	15077.4570 <sup>a</sup>
33	3	31	15077.4570 <sup>a</sup>
33	3	30	15535.5755 <sup>a</sup>
33	4	30	15535.5755 <sup>a</sup>
34	0	34	14617.4624 <sup>a</sup>
34	1	34	14617.4624 <sup>a</sup>
34	1	33	15202.8292 <sup>a</sup>
34	2	33	15202.8292 <sup>a</sup>
34	2	32	15694.4241 <sup>a</sup>
34	3	32	15694.4241 <sup>a</sup>
34	3	31	16163.4249 <sup>a</sup>
34	4	31	16163.4249 <sup>a</sup>
34	4	30	16593.1373 <sup>a</sup>
34	5	30	16593.1373 <sup>a</sup>
35	0	35	15226.7754 <sup>a</sup>
35	1	35	15226.7754 <sup>a</sup>
35	2	34	15827.1128
35	3	32	16804.7609 <sup>a</sup>
35	4	32	16804.7609 <sup>a</sup>
36	0	36	15850.9489
36	1	36	15850.9489
36	1	35	16466.6258
36	2	35	16466.6258

Table 4 (continued)

$J$	$K_a$	$K_c$	Wavenumber
36	2	34	16968.1804
36	3	34	16968.1804
37	0	37	16489.8206
37	1	37	16489.8206
(110)			
28	0	28	12671.4844
28	1	28	12671.4844
28	1	27	13246.5813
28	2	27	13246.5813
28	3	26	13699.4792
29	0	29	13182.9526
29	1	29	13182.9526
29	1	28	13773.7550
29	2	27	14239.5862
30	0	30	13710.5145 <sup>a</sup>
30	1	30	13710.5145 <sup>a</sup>
31	0	31	14254.0956
31	1	31	14254.0956
32	0	32	14813.9119
32	1	32	14813.9119
(011)			
31	0	31	14382.9537 <sup>a</sup>
31	1	31	14382.9537 <sup>a</sup>
31	1	30	15012.3386 <sup>a</sup>
31	2	30	15012.3386 <sup>a</sup>
32	0	32	14944.2833 <sup>a</sup>
32	1	32	14944.2833 <sup>a</sup>
33	0	33	15521.6588 <sup>a</sup>
33	1	33	15521.6588 <sup>a</sup>
33	2	32	16193.0260
34	0	34	16115.0888 <sup>a</sup>
34	1	34	16115.0888 <sup>a</sup>
35	0	35	16724.6749
35	1	35	16724.6749
(120)			
21	0	21	11064.3886 <sup>a</sup>
21	1	21	11064.3131
21	1	20	11549.3003 <sup>a</sup>
21	2	19	11933.5430 <sup>a</sup>
21	3	18	12286.4048 <sup>a</sup>
23	0	23	11868.6405
23	1	22	12404.6686
23	2	21	12812.5262
(021)			
32	0	32	16466.3895 <sup>a</sup>
32	1	32	16466.3895 <sup>a</sup>
(200)			
20	0	20	11097.5337 <sup>a</sup>
20	1	20	11097.5337 <sup>a</sup>
20	1	19	11448.4725 <sup>a</sup>
20	2	19	11448.4725 <sup>a</sup>
20	3	18	11763.8805 <sup>a</sup>
20	4	17	12048.7575 <sup>a</sup>
21	0	21	11473.2241 <sup>a</sup>
21	1	21	11473.2241 <sup>a</sup>
22	0	22	11865.4579
22	1	22	11865.4579
22	1	21	12250.8782
22	2	21	12250.8782
23	0	23	12274.1573 <sup>a</sup>
23	1	23	12274.1573 <sup>a</sup>

<sup>a</sup> Levels labeled have been confirmed by combination differences.

Table 4 (continued)

$J$	$K_a$	$K_c$	Wavenumber
23	1	22	12676.5815 <sup>a</sup>
23	2	22	12676.5815 <sup>a</sup>
23	2	21	13037.8661 <sup>a</sup>
23	3	21	13037.8661 <sup>a</sup>
24	0	24	12699.1791 <sup>a</sup>
24	1	24	12699.1791 <sup>a</sup>
24	1	23	13118.5666 <sup>a</sup>
24	2	23	13118.5666 <sup>a</sup>
24	3	22	13494.2682 <sup>a</sup>
25	0	25	13140.3925 <sup>a</sup>
25	1	25	13140.3925 <sup>a</sup>
25	1	24	13576.6420
25	2	24	13576.6420
26	0	26	13597.6812
26	1	26	13597.6812
27	0	27	14070.9270 <sup>a</sup>
27	1	27	14070.9270 <sup>a</sup>
28	0	28	14559.8711
28	1	28	14559.8711
28	1	27	15046.4377
28	2	27	15046.4377
29	0	29	15064.2789
29	1	29	15064.2789
(101)			
29	0	29	15135.8316
29	1	29	15135.8316
(002)			
26	0	26	13881.7000 <sup>a</sup>
27	0	27	14358.0728 <sup>a</sup>
27	1	27	14358.1455
27	1	26	14821.0409
27	2	26	14821.0409
27	2	25	15233.3060 <sup>a</sup>
27	3	25	15233.3060 <sup>a</sup>
27	3	24	15611.2623
27	4	24	15611.2623
27	4	23	15958.1015 <sup>a</sup>
27	5	22	16276.3080
28	0	28	14850.2901 <sup>a</sup>
28	1	28	14850.2901 <sup>a</sup>
28	1	27	15329.2931 <sup>a</sup>
28	2	27	15329.2931 <sup>a</sup>
28	2	26	15753.5115
28	3	26	15753.5115
28	3	25	16144.2391 <sup>a</sup>
28	4	25	16144.2391 <sup>a</sup>
28	5	24	16503.1292
28	6	23	16832.8937
29	0	29	15358.2833 <sup>a</sup>
29	1	29	15358.2833 <sup>a</sup>
29	1	28	15853.1372 <sup>a</sup>
29	2	28	15853.1372 <sup>a</sup>
29	2	27	16288.6400 <sup>a</sup>
29	3	27	16288.6400 <sup>a</sup>
29	3	26	16691.8220
29	4	26	16691.8220
29	4	25	17062.0055
29	5	24	17402.3850
30	0	30	15881.9026 <sup>a</sup>
30	1	30	15881.9026 <sup>a</sup>
30	1	29	16392.4340
30	2	29	16392.4340
30	2	28	16838.5647
30	3	28	16838.5647

Table 4 (continued)

$J$	$K_a$	$K_c$	Wavenumber
30	3	27	17253.8439 <sup>a</sup>
30	4	27	17253.8439 <sup>a</sup>
30	5	26	17635.1319
30	6	25	17988.9171
31	0	31	16420.9400 <sup>a</sup>
31	1	31	16420.9400 <sup>a</sup>
31	1	30	16947.0228
31	2	30	16947.0228
31	2	29	17403.1731
31	3	29	17403.1731
31	3	28	17830.0626
31	4	28	17830.0626
32	0	32	16975.0995
32	1	32	16975.0995
32	1	31	17516.7366
32	2	31	17516.7366
(031)			
26	0	26	14830.6603 <sup>a</sup>
26	1	26	14829.6624 <sup>a</sup>
26	2	24	15917.6825 <sup>a</sup>
27	1	27	15309.3848 <sup>a</sup>
27	2	26	16026.5598
28	0	28	15820.6570
28	1	27	16557.1781
29	1	29	16352.5276
(012)			
21	0	21	13261.9732 <sup>a</sup>
21	1	21	13261.9732 <sup>a</sup>
21	1	20	13673.3312 <sup>a</sup>
21	2	20	13673.3312 <sup>a</sup>
21	2	19	14035.5136
22	0	22	13652.9553 <sup>a</sup>
22	1	22	13652.9553 <sup>a</sup>
22	2	21	14084.4505 <sup>a</sup>
22	3	20	14459.5103 <sup>a</sup>
22	6	17	15359.8872 <sup>a</sup>
23	0	23	14060.3462 <sup>a</sup>
23	1	23	14060.3462 <sup>a</sup>
23	1	22	14512.0898 <sup>a</sup>
24	1	24	14484.1767 <sup>a</sup>
24	1	23	14956.0772 <sup>a</sup>
24	2	23	14956.0772 <sup>a</sup>
25	0	25	14924.4283 <sup>a</sup>
25	1	25	14924.4283 <sup>a</sup>
26	0	26	15380.3246 <sup>a</sup>
26	1	26	15380.3246 <sup>a</sup>
27	0	27	15852.6949 <sup>a</sup>
27	1	27	15852.6949 <sup>a</sup>
27	1	26	16385.8914 <sup>a</sup>
27	2	26	16385.8914 <sup>a</sup>
29	0	29	16845.3971 <sup>a</sup>
29	1	29	16845.3971 <sup>a</sup>
(121)			
16	0	16	12861.9997 <sup>a</sup>
16	1	16	12861.9970 <sup>a</sup>
16	1	15	13212.5809
16	1	14	13505.7941 <sup>a</sup>
16	2	13	13750.7322 <sup>a</sup>
17	0	17	13165.5455 <sup>a</sup>
17	1	17	13165.5455 <sup>a</sup>
17	2	16	13539.7452 <sup>a</sup>
18	1	17	13883.9766 <sup>a</sup>
18	2	16	14214.0400

(continued on next page)

Table 4 (continued)

$J$	$K_a$	$K_c$	Wavenumber
19	1	19	13822.9574 <sup>a</sup>
19	3	17	14591.0213
19	4	16	14895.0345
20	0	20	14176.8775 <sup>a</sup>
20	1	20	14176.7528 <sup>a</sup>
20	2	19	14623.3093
(022)			
20	0	20	14382.9908 <sup>a</sup>
20	1	19	14826.6092 <sup>a</sup>
20	2	19	14826.6092 <sup>a</sup>
20	4	17	15515.4665 <sup>a</sup>
21	0	21	14754.5857 <sup>a</sup>
21	1	21	14754.6016 <sup>a</sup>
21	2	19	15601.3923 <sup>a</sup>
22	1	22	15142.9020 <sup>a</sup>
22	2	21	15634.2405
22	3	20	16025.9848
23	0	23	15547.2854 <sup>a</sup>
23	1	22	16063.6045 <sup>a</sup>
24	1	24	15968.9740 <sup>a</sup>
24	2	23	16509.3366 <sup>a</sup>
24	3	22	16924.5085 <sup>a</sup>
26	1	26	16861.4715
(131)			
13	1	13	13532.2929
13	3	11	14103.3813
13	4	10	14325.8967
14	0	14	13782.4967
14	1	13	14121.7043
15	0	15	14050.9519
15	1	15	14050.4639 <sup>a</sup>
16	0	16	14336.0995
17	1	17	14638.0818
17	2	16	15060.1423
19	1	19	15292.8586
20	0	20	15646.5620
21	1	21	16016.2291
22	0	22	16403.5502

(200), as well as the (110) level. Additional assignments have also been sunspot spectra covering both the  $L$ - and  $N$ -bands. We are continuing to make progress in assigning the torch spectra.

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### Appendix A. Supplementary data

Supplementary data for this article are available on ScienceDirect ([www.sciencedirect.com](http://www.sciencedirect.com)) and as part of the Ohio State University Molecular Spectroscopy Archives ([http://msa.lib.ohio-state.edu/jmsa\\_hp.htm](http://msa.lib.ohio-state.edu/jmsa_hp.htm)).

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