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Infrared absorption spectra of hot ammonia

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ABSTRACT

Infrared absorption spectra of NH₃ have been obtained at high resolution (0.02 cm⁻¹) at seven temperatures between 296 and 973 K. The spectra were recorded using a Bruker IFS 125 infrared Fourier transform spectrometer in the 2400–5500 cm⁻¹ region and empirical lower state energies have been obtained by comparison of line strengths at different temperatures. Using two reference line lists, quantum number assignments have been made for each temperature for between 1660 and 3020 transitions, with J up to 22. The line lists obtained provide accurate line positions as well as intensities and experimental lower state energies at temperatures relevant for modeling the atmospheres of brown dwarfs and exoplanets.

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1. Introduction

Ammonia (NH₃) is a widely studied molecule in atmospheric science. In the Earth's atmosphere, where sources include animal waste and fertilizers [1], it acts as a precursor in the production of aerosols [2] and particulate matter [3]. Deposition of atmospheric NH₃ can lead to fertilization, a side effect of which may be decreased biodiversity [4].

NH₃ has also been observed in a number of astrophysical environments. In the solar system, it has been detected in the atmospheres of Jupiter [5,6], Saturn [7] and Titan [8] as well as on comets [9,10]. The pure rotational transitions and inversion transitions of NH₃ have been detected in molecular clouds [11,12] making it one of the first extraterrestrial polyatomic molecules discovered. The atmospheres of cool astronomical objects, such as brown dwarfs and exoplanets, are of low enough temperatures to allow small molecules such as NH₃ to form and maintain large enough concentrations to be detected.

Sub-stellar objects with sufficiently low mass (< 0.08 solar masses) are known as brown dwarfs and cannot fuse hydrogen within their cores [13]. The first such object discovered was Gliese 229B [14] and since then a large number of brown dwarfs have been detected with atmospheres cool enough (500 – 2400 K) for a

number of molecular species to exist. Indeed, brown dwarfs, similar to stars, are classified by the presence or absence of particular atomic and molecular features. The hottest such objects, the L dwarfs, contain features from electronic transitions of metal hydrides such as FeH [15] and CrH [16,17]. T dwarfs are distinguished by hot H₂O and CH₄ transitions [18,19]. NH₃ is observed in late T dwarfs, with increasing concentrations as the objects cool [19,20]. As such, NH₃ is expected to characterize a yet cooler class, the Y dwarfs that have an approximate maximum temperature of ~700 K. Y-dwarfs have now been observed and NH₃ appears as a shoulder on an emission feature at 1.58 μm [21,22].

Since the discovery of the first exoplanet in 1995 [23] more than 3500 have been discovered (<http://exoplanet.eu/>). If the planet passes in front of the star it orbits, then it may be detected by observation of periodic decreases in light intensity of the parent star. This technique is known as the transit method and was first used successfully to detect the planet HD 209458b [24] and has since been used by the Kepler mission to observe large numbers of exoplanet candidates [25]. This technique also allows the spectrum of the exoplanet atmosphere to be obtained by recording the transit dips as a function of wavelength [26] and use of the method has resulted in the detection of a number of molecules such as H₂O [27,28], CH₄ [29], CO [30] and CO₂ [31]. As yet, NH₃ has not been observed in the atmosphere of an exoplanet, however a number of modelling studies have predicted its presence in hydrogen-rich hot-Jupiters [32–34].

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NH₃ is a polyatomic molecule having four atoms arranged in a trigonal pyramid structure with C_{3v} symmetry [35]. Ammonia has a complex infrared spectrum with 6 fundamental vibrational modes, two of which are doubly degenerate. The four fundamental vibrational frequencies are: the symmetric stretch at 3336.2 cm⁻¹ (ν_1 , a₁), the symmetric bend at 932.5 cm⁻¹ (ν_2 , a₁), the antisymmetric stretch at 3443.6 cm⁻¹ (ν_3 , e) and the antisymmetric bend at 1626.1 cm⁻¹ (ν_4 , e).

There has been extensive work on the spectroscopy of ammonia in the infrared; for example experimental line lists have been obtained in the 2 μm [36] and 3 μm [37] regions. The most complete line assignments for NH₃ in the infrared are compiled in the HITRAN 2012 database [38]; however HITRAN is intended for applications near room temperature and lacks complete hot band coverage in this region. The hot bands are essential when comparisons are made to high temperature atmospheres. Down *et al.* [39] provide a re-analysis of the HITRAN 2012 NH₃ line list which was used for comparison to HITRAN in this paper; where the original HITRAN 2012 data has been used, the text will mention HITRAN specifically. A comprehensive theoretical line list, BYTe, has been calculated for NH₃ that can be used at temperatures up to 1500 K and contains approximately 1.1 billion transitions [40]. The ro-vibrational calculations from Huang *et al.*, [41,42] and references therein, provide an additional line list for ammonia in this region. High temperature experimental line lists have been obtained using emission spectroscopy at high resolution (0.01 cm⁻¹) from 740–2100 cm⁻¹ [43] and 1650–4000 cm⁻¹ [44]. Absorption spectra at moderate resolution (0.09 cm⁻¹) of hot samples at atmospheric pressure have also been recorded in the 500–2100 cm⁻¹ [45] and 2100–5500 cm⁻¹ [46] spectral regions. Existing experimental line lists have been used in an energy level analysis (MARVEL) to predict a large number of unmeasured lines [47]. Our work reported below details an extension of these experimental line lists, in the 2400–5500 cm⁻¹ region using an improved cell design and technique [48] to record transmission spectra of hot samples. We provide empirical lower state energies, intensities and line positions (± 0.002 cm⁻¹) which can be used in atmospheric models for brown dwarfs and exoplanets.

2. Experimental

Transmission spectra were recorded at seven different temperatures, 23 °C, 200 °C, 300 °C, 400 °C, 500 °C, 600 °C and 700 °C, using a sealed quartz cell, tube furnace and Fourier transform spectrometer. This range includes the effective temperatures of the mid and late T dwarfs [19] and Y dwarfs [22] as well as the coolest observed hot Jupiters [32]. Above 700 °C thermal decomposition of the sample gas reduces the signal to such an extent that reliable spectra could not be obtained.

The transmission spectrum for each temperature is obtained from four individual spectra that, when combined, correct for NH₃ and cell emission; this method has been used previously to record spectra for methane [48]. An absorption spectrum (A_{ab}) is recorded with NH₃ in the cell at temperature with an external emission source (200 W tungsten halogen broadband lamp). A background reference spectrum for the absorption (A_{ref}) is recorded without the NH₃ in the cell. The emission spectrum (B_{em}) is recorded with NH₃ in the cell at temperature without the lamp. The background reference spectrum for emission (B_{ref}) is recorded without the sample in the cell and with the lamp turned off. These spectra are combined to calculate the transmission spectrum for each temperature as

$$\tau = \frac{A_{\text{ab}} - B_{\text{em}}}{A_{\text{ref}} - B_{\text{ref}}}$$

Table 1
Experimental conditions.

Parameter	Value
Spectral region (cm ⁻¹)	2400–5500
Detector	InSb
Beamsplitter	CaF ₂
Spectrometer Windows	CaF ₂
Lens	CaF ₂
Filter	Ge
Scans	300
Resolution (cm ⁻¹)	0.01
NH ₃ Pressure (Torr)	40
Zerofilling factor	x 16

Spectra were recorded in the 2400–5500 cm⁻¹ region, providing overlap with previous experimental work [44]. This previous work includes spectral regions that cover $3\nu_2/\nu_2 + \nu_4$ and $\nu_1/\nu_3/2\nu_4$ bands and the work provided here extends this coverage to include spectral regions that contain the $\nu_1 + \nu_2/\nu_2 + \nu_3$ and $\nu_1 + \nu_4/\nu_3 + \nu_4$ bands as well as associated hot bands.

The 50 cm quartz tube sample cell was used under static conditions to contain the NH₃ sample. The cell is contained within the tube furnace which is heated to the appropriate temperature, which is accurate to within ± 10 °C. The cell was aligned with the entrance aperture of a Bruker IFS 125HR Fourier transform spectrometer and radiation was focused into the spectrometer using a CaF₂ lens. The spectrometer used a CaF₂ beamsplitter and an indium antimonide (InSb) detector. Experimental parameters are summarized in Table 1.

The resulting transmittance spectra contained a number of H₂O lines which were removed manually using the Bruker OPUS software. After combining the spectra at each temperature to calculate a transmission spectrum, peaks were picked using WSpectra [49] to measure their position and intensity. Lines were calibrated for line position and intensity by matching strong and isolated lines that are also found in the HITRAN 2012 database. The number of lines found for each temperature is given in Table 2.

The line intensity equation gives the line intensity as a function of temperature,

$$S' = \frac{2\pi^2\nu_{10}S_{JJ'}}{2\epsilon_0hcQ} \exp\left(-\frac{E''}{kT}\right) \left[1 - \exp\left(-\frac{h\nu_{10}}{kT}\right)\right]$$

where ν_{10} is the line frequency, $S_{JJ'}$ is the line strength, ϵ_0 is the permittivity of free space, h is the Planck constant, c is the speed of light, Q is the internal partition function, E' is the lower state energy, k is the Boltzmann constant and T is the temperature [35]. If a line is measured at different temperatures, the intensity of the line at each temperature may be compared to the intensity at a reference temperature by taking a ratio of the line intensity equation

$$\frac{S'}{S'_0} = \frac{Q_0}{Q} \exp\left(\frac{E''}{kT_0} - \frac{E''}{kT}\right) \left[\frac{1 - \exp\left(-\frac{h\nu_{10}}{kT}\right)}{1 - \exp\left(-\frac{h\nu_{10}}{kT_0}\right)}\right]$$

where S'_0 and T_0 refer to the line intensity and temperature of the reference measurement. For these results 400 °C was used as the reference as this temperature contained the most lines with which to compare the line intensities of other temperatures. If there was no line measured at 400 °C then the reference used was the temperature with the next greatest number of lines. The partition function used for NH₃ was obtained from Yurchenko *et al.* [40], which for our temperature range is essentially the same as the more recent values of Sousa-Silva *et al.* [50].

The lower state energy can be obtained from a rearranged version of the line intensity equation

Table 2
Total number lines, the number of added lines, partition functions [40] and tentative assignments made via comparison with each line list at each temperature. The final line list includes a number of HITRAN lines which replaced saturated lines in the experimental spectra, the number of these is given in the HITRAN additions column.

Temperature (K)	Number of lines	HITRAN additions	Partition function	Down <i>et al.</i> assignments	MARVEL assignments
296	8494	722	1769.563	2136	1813
473	16019	789	3792.440	3020	2721
573	20300	825	5320.817	2976	2839
673	23010	834	7252.482	2913	2906
773	20066	713	9705.249	2569	2594
873	18794	691	12818.840	2296	2346
973	9985	598	16759.410	1660	1178

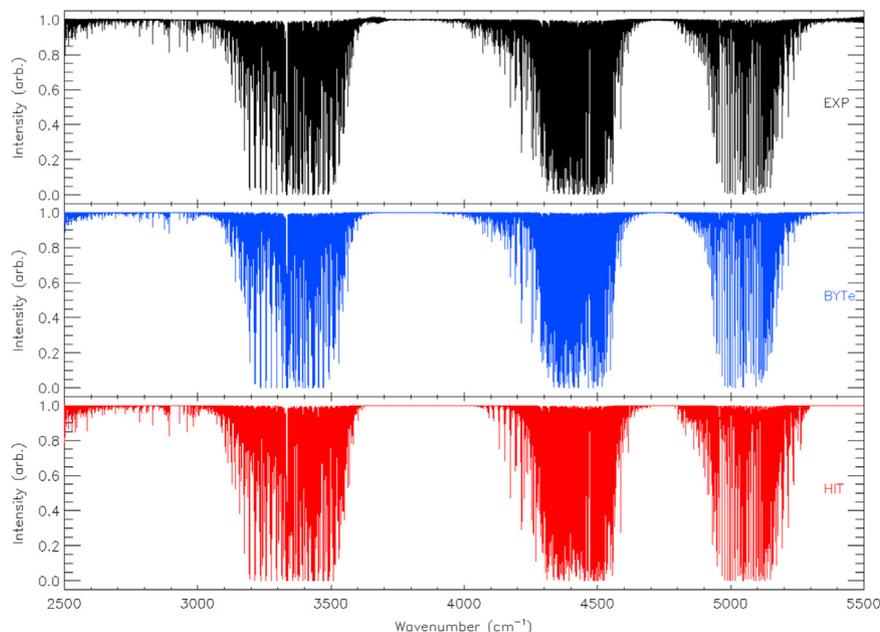


Fig. 1. Overview of spectra of NH₃ showing experimental data (top panel) and simulated spectra created to match the conditions of the experiment from the BYTe [40] (center panel) and Down *et al.* [39] (lower panel) line lists at 296 K.

$$\ln \left(\frac{SQ}{S_0 Q_0} \left[\frac{1 - \exp\left(-\frac{h\nu_{10}}{kT}\right)}{1 - \exp\left(-\frac{h\nu_{10}}{kT_0}\right)} \right] \right) = \frac{E''}{kT_0} - \frac{E''}{kT}$$

where the intensity ratio on the left hand side can be plotted against $1/kT$. Plotted for multiple temperatures, the gradient of this line gives the lower state energy.

3. Results and discussion

Forty Torr of NH₃ was used in the sample cell for these measurements in order to observe weaker transitions. However, at this pressure, many strong lines become saturated, particularly in the Q branch of the ν_1 fundamental mode and the lines are pressure-broadened. The resulting line lists are provided at each temperature (as [Supplementary material](#)) and contain line position, line intensity and lower state energy (if calculated). For observed lines that could be assigned using the MARVEL or Down *et al.* line lists, the corresponding position, intensity, lower state energy and quantum number assignment are provided. Observed saturated lines, for which accurate positions and intensities could not be measured, were excluded and replaced with HITRAN lines with intensities greater than 2.0×10^{-20} cm molecule⁻¹. The experimental positions were calibrated to strong lines in HITRAN, and the mean difference between all matched lines is approximately 4.0×10^{-4} cm⁻¹ with a standard deviation of 3.0×10^{-4} cm⁻¹ after

calibration. Similarly, the accuracy of the intensities of the experimentally obtained lines is measured by comparison with matched HITRAN lines; for these the mean difference is within 20%.

The top panels of [Figs. 1](#) and [2](#) show overviews of the recorded spectra at 296 and 773 K, respectively. The middle panels show simulated spectra from the BYTe line list and the bottom panels similarly show simulated spectra from the Down *et al.* line list. The spectral simulation uses the Reference Forward Model [51], a line-by-line radiative transfer model, to calculate the transmission through a 50 cm cell with the experimental pressure and temperature. Individual line broadening parameters for the lines in BYTe were obtained as averages from HITRAN 2012 and applied to all lines. The Down *et al.* spectral simulation uses the line positions and intensities from HITRAN 2012 as well as the averaged broadening parameters applied globally, therefore allowing for comparisons between Down *et al.*, BYTe and our measurements.

The experimental spectrum at 296 K ([Fig. 1](#)) shows some saturated lines of the ν_1 fundamental, $\nu_1 + \nu_2/\nu_2 + \nu_3$ and $\nu_1 + \nu_4/\nu_3 + \nu_4$ modes which have been replaced in the final line list. In general, there is good agreement between the three spectra at 296 K. For 773 K, the experimental and BYTe spectra in [Fig. 2](#) clearly show a number of additional hot bands that are not present in the Down *et al.* spectrum, most notably the Q-branch of the $\nu_1 + 3\nu_3 - 3\nu_3$ band at around 3850 cm⁻¹ and the $\nu_2 + \nu_3 - \nu_2$ band from 4000 cm⁻¹ to 4200 cm⁻¹. Other hot bands are included in the 3000–3600 cm⁻¹ regions and 4200–4700 cm⁻¹ regions although they are not visible because of the density of lines. Some

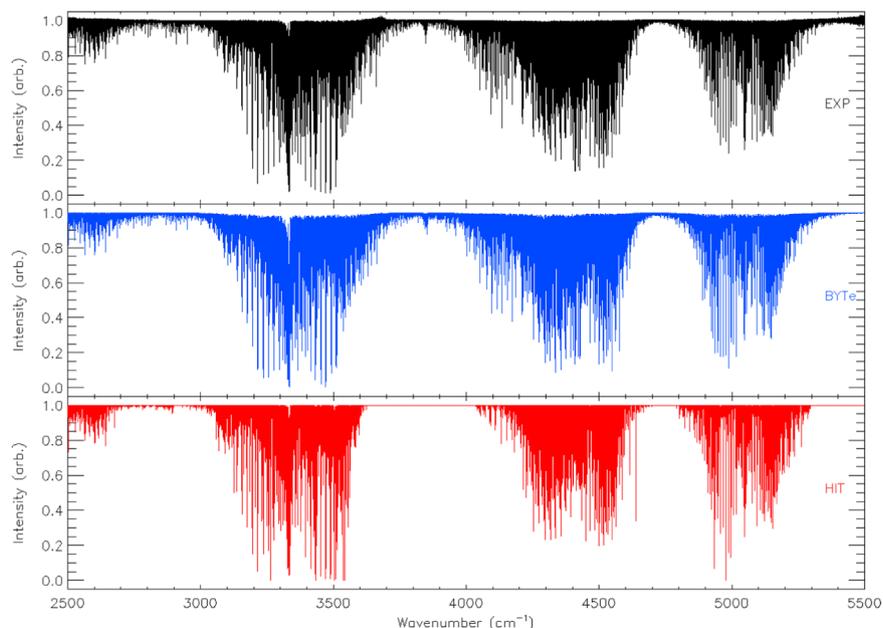


Fig. 2. Overview of spectra of NH_3 showing experimental data (top panel) and simulated spectra created to match the conditions of the experiment from the BYTe [40] (center panel) and Down et al. [39] (lower panel) line lists at 773 K.

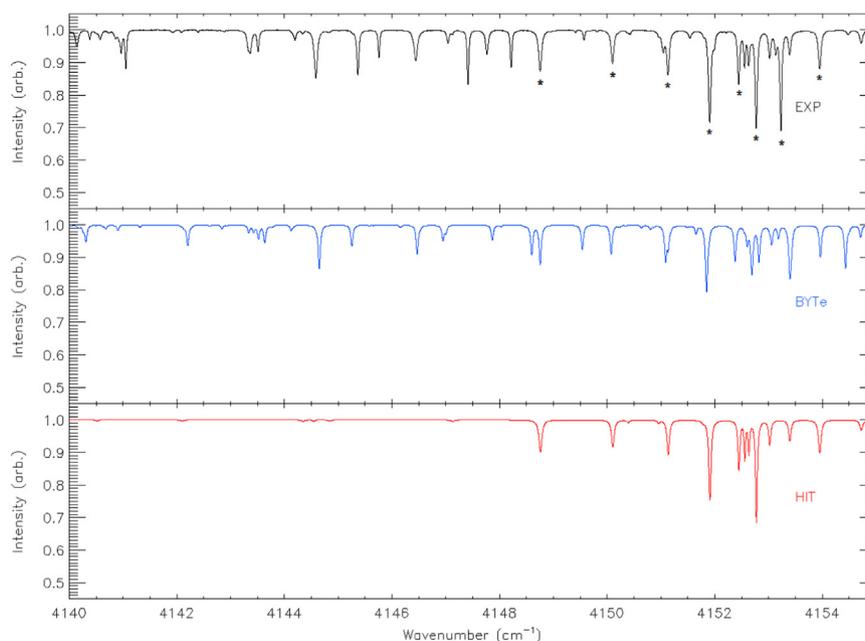


Fig. 3. Detailed spectra of NH_3 showing experimental data (top panel) and simulated spectra created to match the conditions of the experiment from the BYTe [40] (center panel) and Down et al. [39] (lower panel) line lists at 296 K. Experimental lines that have been assigned by matching are marked by asterisks.

lines in HITRAN have an incorrect lower state energy which can be seen clearly in Fig. 2 at, for example, around 4150 cm^{-1} where an error in lower state energy results in a ‘bad’ extrapolation with temperature.

A more detailed view of short segments of the three spectra is shown in Figs. 3 and 4. Fig. 3 again shows the three spectra for 296 K. There are a number of lines that are assignable in both the experimental and BYTe spectra; one example being the strong line at 4151.8 cm^{-1} (the R(6) line of the $\nu_1 + \nu_2$ band) that appears in all three spectra within 0.1 cm^{-1} with a similar intensity. Even at room temperature, however, there are a number of lines in the experimental spectrum that are missing from Down *et al.* and a number of inconsistencies between the experimental spectrum and the BYTe simulation. Lines in the experimental line lists which

have been matched to lines in either reference line list are marked by an asterisk. At 773 K , the hot lines in the experimental spectrum at 4141.0 cm^{-1} ($0\ 2\ 1^1\ 0^0 - 0\ 1\ 0^0\ 0^0\ \text{P}(5\ 4)^+$, notation in the form of *excited state* $\nu_1\ \nu_2\ \nu_3^l\ \nu_4^l - \text{lower state } \nu_1\ \nu_2\ \nu_3^l\ \nu_4^l\ \text{P}(J^* K')^{+/-}$ where $+/-$ is the lower state parity [40]; the inversional parity for all three lines is given as 0 in the BYTe line list), 4147.4 cm^{-1} ($0\ 2\ 1^1\ 0^0 - 0\ 1\ 0^0\ 0^0\ \text{P}(5\ 5)^-$) and 4153.2 cm^{-1} ($0\ 2\ 1^1\ 0^0 - 0\ 1\ 0^0\ 0^0\ \text{P}(3\ 3)^-$) can be matched to lines in the BYTe spectrum, although their positions are shifted. In this region the shift is over 1 cm^{-1} towards higher wavenumbers, although in other regions the shift value and direction vary. While the strong lines can be identified, it becomes more problematic to assign weaker features from BYTe primarily due to the observed shifts. The lines in Fig. 4 which are not matched (i.e., those without an asterisk)

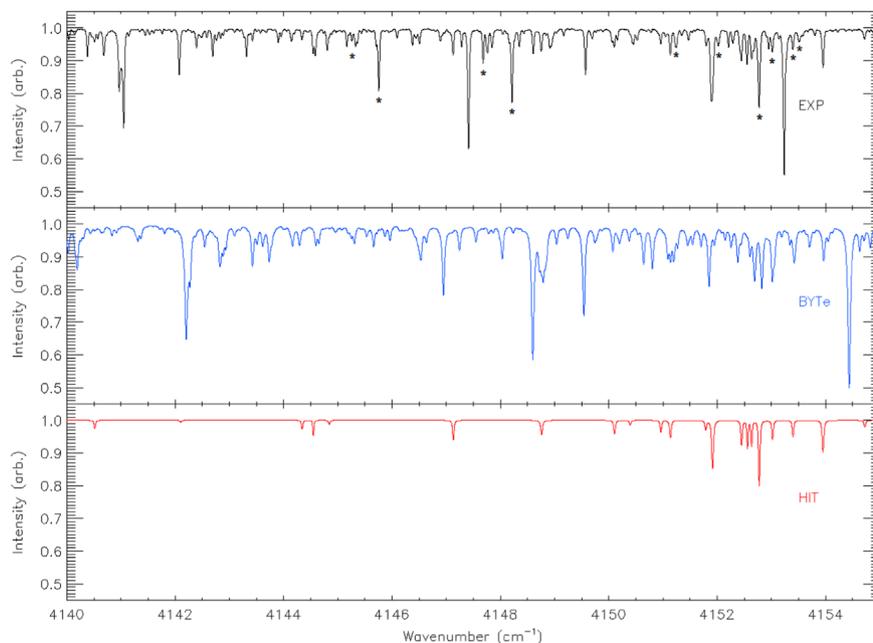


Fig. 4. Detailed spectra of NH_3 showing experimental data (top panel) and simulated spectra created to match the conditions of the experiment from the BYTe [40] (center panel) and Down et al. [39] (lower panel) line lists at 773 K. Experimental lines that have been assigned by matching are marked by asterisks.

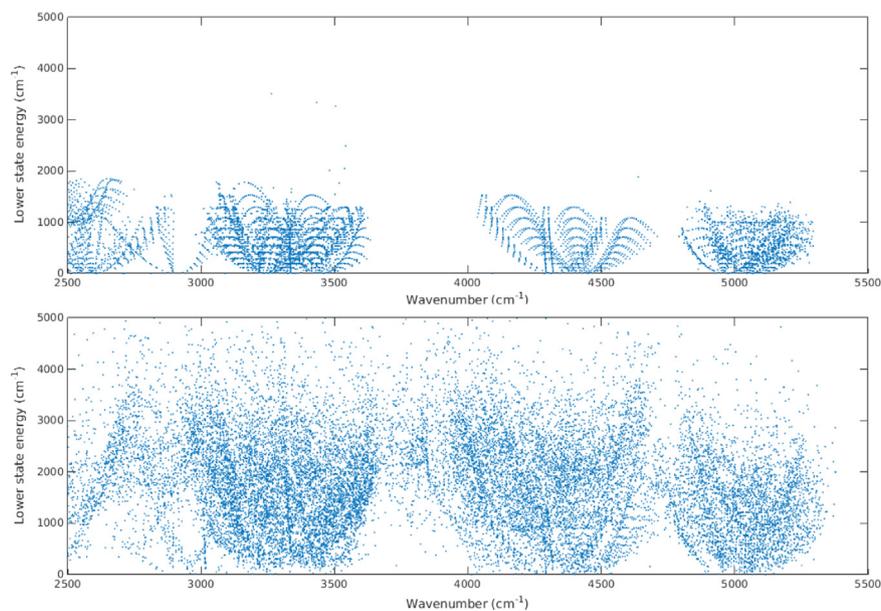


Fig. 5. Comparison of lower state energies from the HITRAN line list [38] (top panel) and from the experimental line list (lower panel) as a function of line position.

include a number of strong lines, and these are mostly hot lines. This demonstrates the difficulty in assigning lines in regions with overlapping hot bands where the line position from calculated spectra is of insufficient accuracy.

The empirical lower state energies are plotted against line position in Fig. 5 in the lower panel along with those from Down *et al.* in the upper panel. The near-vertical lines belong to Q branches and the parabolic features to the left and right of the Q-branches are from the P and R branches, respectively. The accuracy of these empirical lower state energies is difficult to quantify, but has been discussed in our earlier work [43,44] and is dependent on a number of factors, such as the line strength, blending effects, and the number of temperatures the line is measured. For strong, isolated lines, E' is typically within 10% of ab initio E' calculations, however for weaker lines, the differences can be larger. Care should therefore be taken when using these E' values

since they are not necessarily sufficient for spectroscopic assignments but when used as a whole, they are appropriate for extrapolating the NH_3 spectrum to nearby temperatures when no calculated E' is available. Fig. 6 displays the empirical E' plotted against corresponding values from Down *et al.* [39]. The expected linear ratio is also plotted and it can be seen that the majority of points lie close to this line, with a small number of points displaying a larger difference. The gradient of the data in Fig. 6 may be slightly higher than unity and the intercept x-intercept appears to be above zero. These are most likely due to small errors in calibration which is made difficult by slight sample decomposition at higher temperatures.

There are a number of notable differences between the Down *et al.* lower state energies and those obtained by experiment (Fig. 5). The incompleteness of Down *et al.* (and therefore HITRAN) can be seen between 3600 and 4100 cm^{-1} ; experimental data has

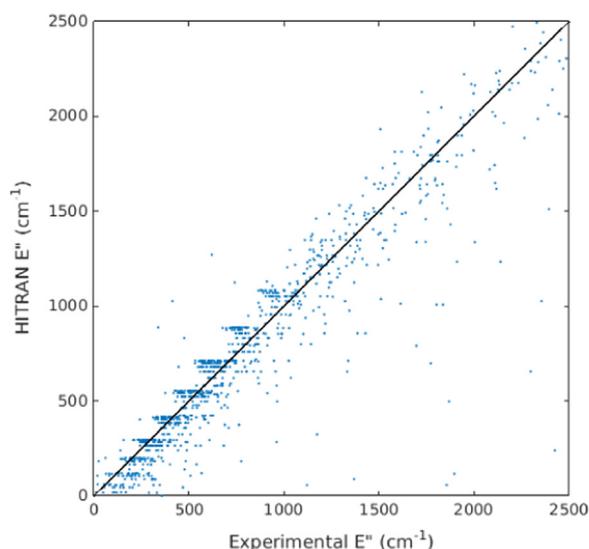


Fig. 6. Comparison of matched lower state energies obtained from experiment (x-axis) and Down et al. [39] (y-axis).

now been obtained in this region. There are also very few lower state energy values listed in HITRAN with energies above 2000 cm^{-1} ; some of these are from fundamental and combination transitions and can be seen in the bottom panel, particularly the Q branch of the ν_1 band and the P and R branches of the $\nu_1 + \nu_2$ band. However, most of the lines with high lower state energies are from hot transitions. The most notable 'hot' band in the lower panel is the strong Q branch at a position 3850 cm^{-1} , extending from a lower state energy of $2000\text{--}3000\text{ cm}^{-1}$, which is not provided in HITRAN. Fig. 2 This feature can be seen in both the experimental data (top) and BYTE (center) panels but not in the Down *et al.* (lower) panel of Fig. 2.

The MARVEL analysis of NH_3 [45,46] and Down et al. NH_3 database [39] are line lists in the region of interest which contain quantum number assignments. These recent analyses have used experimental line positions to extend quantum number assignments. MARVEL provides a self-consistent list of fully assigned line positions, but no associated line intensities. Therefore, rather than using MARVEL directly, the MARVEL adjusted BYTE line list BARVEL is used, which was computed by replacing BYTE energies with the corresponding values from MARVEL. Disregarding the 13% of MARVEL states that could not be matched to corresponding BYTE states, BARVEL provides intensity data for roughly 18,500 (66%) of MARVEL transitions, plus many predicted transitions between MARVEL energy levels that have not been observed experimentally. The remaining line list consists of transitions for which either one or both of the states takes its energy from BYTE. Only the transitions for which both upper and lower state had their values taken from MARVEL were used to avoid any discrepancies in line positions due to the inaccuracy of BYTE.

The positions and intensities from both the BARVEL and Down et al. lists may be compared to those obtained experimentally to provide tentative assignments. Lines with positions within 0.002 cm^{-1} and intensities between 0.5 and 2.0 times that of the experimental lines were considered matched, with the number of assigned lines from each reference list given in Table 2. It should be noted that many of these assignments are for the same lines at different temperatures. Quantum number assignments can therefore be provided for this experimental line list. The assignments from both BARVEL and Down *et al.* are included in the final line list for. In general, these comparisons are used only to provide quantum numbers for completeness and are not used as a measure of the quality of our new list.

4. Conclusions

High resolution transmission spectra of ammonia recorded at seven temperatures between $22\text{ }^\circ\text{C}$ and $700\text{ }^\circ\text{C}$ were used to calculate lower state energies in the $2400\text{--}5500\text{ cm}^{-1}$ region. These spectra include several hot bands which are not present in HITRAN even at room temperature with tentative assignments of J up to 22. Therefore, this work currently provides the most accurate line list for simulating room temperature (and higher) NH_3 observations in this spectral region. Extrapolating room temperature HITRAN results to higher temperatures results in a substantial decrease in total intensity as the hot bands are not included (Fig. 4) and this work provides intensity estimates for the hot bands in this region.

Calculated line lists such as BYTE [40] provide more lines than are experimentally observable due to the weakness of millions of lines in the spectrum and the sensitivity required to measure them. Calculations give a much more complete total intensity sum over multiple bands, however, positions of individual lines in BYTE are less well determined than in the experiment measurements.

The experimental line lists obtained can be used directly in atmospheric models for brown dwarfs and exoplanets. The line lists can also be used to create template spectra for cross correlation with high resolution spectra of planetary systems. The cross correlation method for detecting molecules on exoplanets, such as those used to detect CO on HD 209458b [52] by analyzing Doppler-shifted spectral features of a transiting exoplanet, require high resolution template spectra with accurate line positions and intensities. The line lists of this work satisfy both the high temperature and high resolution requirements for current exoplanetary research.

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Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at <http://dx.doi.org/10.1016/j.jqsrt.2017.02.012>.

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