# The high resolution infrared spectroscopy of cyanogen di- N -oxide (ONCCNO) 

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The high-resolution infrared absorption spectrum of the oxalodinitrile di-N-oxide (ONCCNO) molecule has been recorded in the gas phase with a Fourier transform spectrometer at a resolution of $0.003 \mathrm{~cm}^{-1}$. No previous high-resolution spectra have been recorded for this semistable palindromic molecule. On the basis of the $2: 1$ intensity alternation in the rotational lines caused by nitrogen nuclear spin statistics, the ONCCNO molecule appears to be linear. A quasilinear structure, however, cannot be ruled out at this stage of the analysis. The $\nu_{4}$ and $\nu_{5}$ fundamental modes at $2246.04055(23) \mathrm{cm}^{-1}$ and $1258.47530(11) \mathrm{cm}^{-1}$ have been analyzed to give ground state rotational constants of $B_{0}=0.04220210(96) \mathrm{cm}^{-1}$ and $D_{0}=8.77(70) \times 10^{-10} \mathrm{~cm}^{-1}$. By fixing the CN and NO bond lengths to 1.1923 and $1.1730 \AA$, respectively, the $\mathrm{C}-\mathrm{C}$ bond length was determined to be $1.3329 \AA$ using the $B_{0}$ value. This short $\mathrm{C}-\mathrm{C}$ bond length is thus similar to that observed for a carbon-carbon double bond. © 1995 American Institute of Physics.

## INTRODUCTION

The cyanogen di-N-oxide (oxalodinitrile di-N-oxide, ONCCNO) molecule was first prepared at the beginning of this century. ${ }^{1}$ However, its chemical formula and isolation were not achieved until the early 1960s by Grundmann. ${ }^{2,3}$ The ONCCNO molecule was prepared in organic solutions by HCl elimination from the stable dichloroglyoxime $(\mathrm{HON}=\mathrm{C}(\mathrm{Cl})-(\mathrm{Cl}) \mathrm{C}=\mathrm{NOH})$ precursor. Two strong infrared absorption peaks at 2190 and $1235 \mathrm{~cm}^{-1}$ were observed in $\mathrm{CCl}_{4}$ solution. The ultraviolet spectrum, which contains maxima at 312,295 , and 262 nm , can be obtained with a


FIG. 1. (a) An overview of the $\nu_{4} \sigma_{u} \mathrm{CNO}$ antisymmetric stretching mode of the ONCCNO molecule. Note the lines due to the HNCO molecule near $2270 \mathrm{~cm}^{-1}$. (b) An expanded portion of the $R$ branch $2: 1$ intensity alternation caused by the nuclear spin statistics of two equivalent nitrogen nuclei in a linear molecule.
freshly prepared solution of ONCCNO in $n$-hexane before polymerization occurs to make polyfuroxan. ${ }^{3}$ Since dilute solutions are stable at $0^{\circ} \mathrm{C}$ for several hours, ONCCNO has been widely used in organic chemistry for 1,3-dipolar cycloaddition reactions. ${ }^{4}$

The ONCCNO molecule is a candidate for astrophysical observation since it contains only the relatively abundant elements $\mathrm{C}, \mathrm{N}$, and O . The symmetric CNO dimer structure is also of spectroscopic and structural interest since there is the possibility of quasilinear behavior as found in HCNO. ${ }^{5}$ The symmetric linear CNO dimer ONCCNO has no dipole moment so that pure rotational spectra will be very weak and infrared or ultraviolet observations will be necessary.


FIG. 2. Portions of $R$ and $P$ branches of the $\nu_{4}$ mode of the ONCCNO molecule. The $J$ quantum numbers are for the fundamental band and the perturbed $R(96)$ and $P(98)$ lines are marked by arrows.

TABLE I. The line list for the $\nu_{4}$ vibrational mode of ONCCNO $\left(\mathrm{cm}^{-1}\right)$.

| Line | OBS | $\begin{gathered} \mathrm{OMC}^{\mathrm{a}} \\ \times 10^{4} \end{gathered}$ | Line | OBS | $\begin{gathered} \mathrm{OMC} \\ \times 10^{4} \end{gathered}$ | Line | OBS | $\begin{gathered} \mathrm{OMC} \\ \times 10^{4} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $P(101)$ | 2236.5060 | -8 | $P(53)$ | 2241.3072 | 27 | $P(5)$ | 2245.6118 | -49 |
| $P(100)$ | 2236.6118 | -1 | $P(52)$ | 2241.4013 | 22 | $P(4)$ | 2245.6935 | -83 |
| $P(99)$ | 2236.7160 | -9 | $P(51)$ | 2241.4966 | 32 | $P(3)$ | 2245.7812 | -56 |
| $P(98)$ | 2236.8269 | 53 | $P(50)$ | 2241.5893 | 17 | $P(2)$ | 2245.8638 | -77 |
| $P(97)$ | 2236.9284 | 23 | $P(49)$ | 2241.6858 | 43 | $R(1)$ | 2246.2950 | -54 |
| $P(96)$ | 2237.0320 | 17 | $P(48)$ | 2241.7782 | 30 | $R(2)$ | 2246.2950 | 24 |
| $P(95)$ | 2237.1364 | 19 | $P(47)$ | 2241.8715 | 28 | $R(3)$ | 2246.3702 | -61 |
| $P(94)$ | 2237.2421 | 39 | $P(46)$ | 2241.9647 | 27 | $R(4)$ | 2246.4529 | -69 |
| $P(93)$ | 2237.3436 | 17 | $P(45)$ | 2242.0581 | 30 | $R(5)$ | 2246.5382 | -49 |
| $P(92)$ | 2237.4446 | -8 | $P(44)$ | 2242.1510 | 31 | $R(6)$ | 2246.6200 | -61 |
| $P(91)$ | 2237.5485 | 0 | $P(43)$ | 2242.2427 | 21 | $R(7)$ | 2246.7011 | -80 |
| $P(90)$ | 2237.6514 | -1 | $P(42)$ | 2242.3361 | 30 | $R(8)$ | 2246.7855 | -63 |
| $P(89)$ | 2237.7516 | -26 | $P(41)$ | 2242.4287 | 33 | $R(9)$ | 2246.8691 | -53 |
| $P(88)$ | 2237.8557 | -11 | $P(40)$ | 2242.5205 | 31 | $R(10)$ | 2246.9513 | -54 |
| $P(87)$ | 2237.9575 | -16 | $P(39)$ | 2242.6120 | 27 | $R(11)$ | 2247.0333 | -56 |
| $P(86)$ | 2238.0600 | -12 | $P(38)$ | 2242.7041 | 31 | $R(12)$ | 2247.1181 | -28 |
| $P(85)$ | 2238.1623 | -7 | $P(37)$ | 2242.7953 | 29 | $R(13)$ | 2247.1974 | -53 |
| $P(84)$ | 2238.2651 | 4 | $P(36)$ | 2242.8852 | 16 | $R(14)$ | 2247.2806 | -37 |
| $P(83)$ | 2238.3634 | -27 | $P(35)$ | 2242.9770 | 23 | $R(15)$ | 2247.3624 | -33 |
| $P(82)$ | 2238.4667 | -6 | $P(34)$ | 2243.0666 | 10 | $R(16)$ | 2247.4447 | -22 |
| $P(81)$ | 2238.5661 | -22 | $P(33)$ | 2243.1527 | -36 | $R(17)$ | 2247.5253 | -27 |
| $P(80)$ | 2238.6689 | -1 | $P(32)$ | 2243.2526 | 59 | $R(18)$ | 22247.6062 | -26 |
| $P(79)$ | 2238.7669 | -26 | $P(31)$ | 2243.3397 | 27 | $R(19)$ | 2247.6881 | -14 |
| $P(78)$ | 2238.8686 | -12 | $P(30)$ | 2243.4298 | 28 | $R(20)$ | 2247.7683 | -17 |
| $P(77)$ | 2238.9684 | -15 | $P(29)$ | 2243.5196 | 27 | $R(21)$ | 2247.8501 | -1 |
| $P(76)$ | 2239.0690 | -7 | $P(28)$ | 2243.6081 | 15 | $R(22)$ | 2247.9302 | -1 |
| $P(75)$ | 2239.1677 | -16 | $P(27)$ | 2243.6979 | 18 | $R(23)$ | 2248.0123 | 21 |
| $P(74)$ | 2239.2681 | -6 | $P(26)$ | 2243.7865 | 11 | $R(24)$ | 2248.0898 | -1 |
| $P(73)$ | 2239.3673 | -5 | $P(25)$ | 2243.8740 | -5 | $R(25)$ | 2248.1665 | -29 |
| $P(72)$ | 2239.4666 | -1 | $P(24)$ | 2243.9639 | 5 | $R(26)$ | 2248.2514 | 27 |
| $P(71)$ | 2239.5654 | 0 | $P(23)$ | 2244.0515 | -7 | $R(27)$ | 2248.3291 | 13 |
| $P(70)$ | 2239.6629 | -10 | $P(22)$ | 2244.1396 | -10 | $R(28)$ | 2248.4094 | 27 |
| $P(69)$ | 2239.7629 | 7 | $P(21)$ | 2244.2282 | -8 | $R(29)$ | 2248.4895 | 41 |
| $P(68)$ | 2239.8599 | -3 | $P(20)$ | 2244.3156 | -15 | $R(30)$ | 2248.5708 | 68 |
| $P(67)$ | 2239.9585 | 4 | $P(19)$ | 2244.4029 | -22 | $R(31)$ | 2248.6387 | -36 |
| $P(66)$ | 2240.0548 | -8 | $P(18)$ | 2244.4907 | -22 | $R(32)$ | 2248.7207 | 3 |
| $P(65)$ | 2240.1539 | 8 | $P(17)$ | 2244.5765 | -39 | $R(33)$ | 2248.7988 | 5 |
| $P(64)$ | 2240.2507 | 5 | $P(16)$ | 2244.6640 | -38 | $R(34)$ | 2248.8785 | 25 |
| $P(63)$ | 2240.3482 | 11 | $P(15)$ | 2244.7505 | -45 | $R(35)$ | 2248.9567 | 31 |
| $P(62)$ | 2240.4448 | 9 | $P(14)$ | 2244.8374 | -46 | $R(36)$ | 2249.0333 | 24 |
| $P(61)$ | 2240.5415 | 11 | $P(13)$ | 2244.9236 | -52 | $R(37)$ | 2249.1097 | 17 |
| $P(60)$ | 2240.6367 | 1 | $P(12)$ | 2245.0107 | -48 | $R(38)$ | 2249.1880 | 30 |
| $P(59)$ | 2240.7338 | 11 | $P(11)$ | 2245.0969 | -50 | $R(39)$ | 2249.2650 | 33 |
| $P(58)$ | 2240.8287 | 2 | $P(10)$ | 2245.1795 | -87 | $R(40)$ | 2249.3414 | 33 |
| $P(57)$ | 2240.9276 | 35 | $P(9)$ | 2245.2674 | -68 | $R(41)$ | 2249.4181 | 36 |
| $P(56)$ | 2241.0207 | 11 | $P(8)$ | 2245.3533 | -68 | $R(42)$ | 2249.4935 | 29 |
| $P(55)$ | 2241.1168 | 21 | $P(7)$ | 2245.4385 | -73 | $R(43)$ | 2249.5689 | 25 |
| $P(54)$ | 2241.2121 | 23 | $P(6)$ | 2245.5231 | -83 | $R(44)$ | 2249.6450 | 29 |
| $R(45)$ | 2249.7211 | 35 | $R(73)$ | 2251.7420 | -6 | $R(101)$ | 2253.5926 | 14 |
| $R(46)$ | 2249.7967 | 38 | $R(74)$ | 2251.8111 | -6 | $R(102)$ | 2253.6543 | 3 |
| $R(47)$ | 2249.8708 | 29 | $R(75)$ | 2251.8792 | -13 | $R(103)$ | 2253.7177 | 11 |
| $R(48)$ | 2249.9443 | 16 | $R(76)$ | 2251.9475 | -16 | $R(104)$ | 2253.7804 | 14 |
| $R(49)$ | 2250.0198 | 24 | $R(77)$ | 2252.0167 | -9 | $R(105)$ | 2253.8425 | 13 |
| $R(50)$ | 2250.0941 | 24 | $R(78)$ | 2252.0846 | -11 | $R(106)$ | 2253.9047 | 15 |
| $R(51)$ | 2250.1678 | 18 | $R(79)$ | 2252.1528 | -8 | $R(107)$ | 2253.9687 | 37 |
| $R(52)$ | 2250.2425 | 26 | $R(80)$ | 2252.2202 | -11 | $R(108)$ | 2254.0281 | 15 |
| $R(53)$ | 2250.3149 | 12 | $R(81)$ | 2252.2873 | -15 | $R(109)$ | 2254.0896 | 16 |
| $R(54)$ | 2250.3890 | 18 | $R(82)$ | 2252.3550 | -11 | $R(110)$ | 2254.1513 | 21 |
| $R(55)$ | 2250.4622 | 16 | $R(83)$ | 2252.4212 | -19 | $R(111)$ | 2254.2115 | 12 |
| $R(56)$ | 2250.5334 | -3 | $R(84)$ | 2252.4884 | -15 | $R(112)$ | 2254.2727 | 16 |
| $R(57)$ | 2250.6075 | 9 | $R(85)$ | 2252.5563 | -2 | $R(113)$ | 2254.3330 | 12 |
| $R(58)$ | 2250.6802 | 10 | $R(86)$ | 2252.6224 | -4 | $R(114)$ | 2254.3936 | 14 |
| $R(59)$ | 2250.7532 | 15 | $R(87)$ | 2252.6876 | -14 | $R(115)$ | 2254.4542 | 17 |
| $R(60)$ | 2250.8248 | 9 | $R(88)$ | 2252.7542 | -6 | $R(116)$ | 2254.5137 | 10 |
| $R(61)$ | 2250.8970 | 10 | $R(89)$ | 2252.8202 | -3 | $R(117)$ | 2254.5730 | 4 |
| $R(62)$ | 2250.9681 | 3 | $R(90)$ | 2252.8850 | -9 | $R(118)$ | 2254.6332 | 8 |

TABLE I. (Continued.)

| Line | OBS | $\begin{gathered} \mathrm{OMC}^{\mathrm{a}} \\ \times 10^{4} \end{gathered}$ | Line | OBS | $\begin{gathered} \mathrm{OMC} \\ \times 10^{4} \end{gathered}$ | Line | OBS | $\begin{gathered} \mathrm{OMC} \\ \times 10^{4} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $R(63)$ | 2251.0405 | 12 | $R(91)$ | 2252.9512 | 0 | $R(119)$ | 2254.6923 | 3 |
| $R(64)$ | 2251.1115 | 9 | $R(92)$ | 2253.0164 | 2 | $R(120)$ | 2254.7497 | -19 |
| $R(65)$ | 2251.1820 | 2 | $R(93)$ | 2253.0807 | -2 | $R(121)$ | 2254.8104 | -4 |
| $R(66)$ | 2251.2515 | -12 | $R(94)$ | 2253.1465 | 10 | $R(122)$ | 2254.8690 | -10 |
| $R(67)$ | 2251.3233 | 0 | $R(95)$ | 2253.2110 | 13 | $R(123)$ | 2254.9284 | -6 |
| $R(68)$ | 2251.3940 | 3 | $R(96)$ | 2253.2800 | 61 | $R(124)$ | 2254.9856 | -23 |
| $R(69)$ | 2251.4636 | -4 | $R$ (97) | 2253.3339 | -39 |  |  |  |
| $R(70)$ | 2251.5333 | -7 | $R(98)$ | 2253.4004 | -10 |  |  |  |
| $R(71)$ | 2251.6030 | -8 | $R(99)$ | 2253.4645 | -4 |  |  |  |
| $R(72)$ | 2251.6724 | -9 | $R(100)$ | 2253.5281 | 0 |  |  |  |

${ }^{\text {a }}$ Observed minus calculated line positions.

Pure samples of ONCCNO are difficult to handle. Crystalline, monomeric ONCCNO is stable at $-78^{\circ} \mathrm{C}$ but begins to visibly decompose near $-45^{\circ} \mathrm{C}$ and explodes a few minutes later at that temperature. ${ }^{6}$ Maier and Teles ${ }^{7}$ have reported the formation of ONCCNO by flash vacuum pyrolysis of dichloroglyoxime followed by condensation of the pyrolysis products diluted with argon on a cold ( 10 K ) window, although they published no spectroscopic data. Until very recently no gas phase measurements have been made. Pasinszki and Westwood ${ }^{8}$ have successfully studied gaseous ONCCNO by He I photoelectron spectroscopy, photoionization mass spectrometry, low-resolution midinfrared spectroscopy as well as $a b$ initio calculations. The strong IR absorption bands at 2226 and $1260 \mathrm{~cm}^{-1}$ correspond to the antisymmetric $\left(\nu_{4} \sigma_{u}\right)$ and symmetric $\left(\nu_{5} \sigma_{u}\right)$ CNO group stretching vibrations, respectively, which were reported earlier at 2190 $\mathrm{cm}^{-1}$ and $1235 \mathrm{~cm}^{-1}$ in $\mathrm{CCl}_{4}$ solution. ${ }^{2}$ Unfortunately, the $a b$ initio calculations did not provide a solid conclusion as to whether the ONCCNO molecule was linear or bent although a linear or quasilinear structure was preferred. Highresolution gas phase spectroscopy can supply additional information on the molecular geometry.

In this work, we report on our high-resolution Fourier transform infrared spectra of the cyanogen di-N-oxide molecule. The $\nu_{4}$ and $\nu_{5}$ vibrational modes have been recorded in the gas phase and these two fundamental bands have been rotationally analyzed.

## EXPERIMENT

The ONCCNO molecule was generated in situ using the same method as described by Pasinszki and Westwood. ${ }^{8}$ Briefly, the thermolysis of dichloroglyoxime in a quartz tube ( 8 mm i.d. by 15 cm ) heated to $550^{\circ} \mathrm{C}$ gives a good yield of ONCCNO plus HCl with only trace amounts of the side products, $\mathrm{NO}, \mathrm{CO}, \mathrm{CO}_{2}$, and HNCO . The experimental setup is typical for absorption spectroscopic work using a cell and a glower external to the spectrometer. The infrared glower was collimated by a parabolic mirror and passed through a 20 cm long absorption cell equipped with KBr windows and entered the spectrometer through the emission port. The thermolysis products were pumped slowly through the gas cell at a pressure of about 250 mTorr .

The high-resolution absorption spectra were recorded with a Bruker FTS 120 HR spectrometer at the University of Waterloo with a resolution of $0.003 \mathrm{~cm}^{-1}$. The $\nu_{4}$ mode, near $2250 \mathrm{~cm}^{-1}$, was recorded with an InSb detector and 46 scans were coadded in the $1800-2900 \mathrm{~cm}^{-1}$ region. A redpass filter with a cutoff at $2900 \mathrm{~cm}^{-1}$ set the upper wave number limit while the lower wave number limit was set by the band gap of the InSb detector. A HgCdTe detector was used to record the $\nu_{5}$ mode at $1260 \mathrm{~cm}^{-1}$ and the spectrum was obtained by coadding 55 scans. Another redpass filter (cutoff at $1672 \mathrm{~cm}^{-1}$ ) set the high wave number limit and the HgCdTe detector response set the lower limit. A KBr beamsplitter was used for both spectra.

## ANALYSIS

The spectral analysis program PC-DECOMP, developed by J. W. Brault, was used for the spectral line measurements. Using this program, the line profiles were fitted with Voigt lineshape functions. The signal-to-noise ratio for the strongest lines in the spectrum was about 5:1 and the precision of the line position measurement is better than $\pm 0.0006 \mathrm{~cm}^{-1}$


FIG. 3. (a) An overview of the $\nu_{5} \sigma_{u}$ CNO symmetric stretching mode of the ONCCNO molecule. (b) An expanded portion of the $P$ branch of the $\nu_{5}$ mode. The marked lines are for the fundamental band and show the $2: 1$ intensity alternation caused by nuclear spin statistics.

TABLE II. The line list for the $\nu_{5}$ vibrational mode of ONCCNO $\left(\mathrm{cm}^{-1}\right)$.

| Line | OBS | $\begin{gathered} \mathrm{OMC}^{\mathrm{a}} \\ \times 10^{4} \end{gathered}$ | Line | OBS | $\begin{gathered} \mathrm{OMC} \\ \times 10^{4} \end{gathered}$ | Line | OBS | $\begin{gathered} \mathrm{OMC} \\ \times 10^{4} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $P(141)$ | 1245.2485 | -8 | $P(91)$ | 1250.2297 | -3 | $P(43)$ | 1254.7237 | 17 |
| $P(140)$ | 1245.3501 | -8 | $P(90)$ | 1250.3266 | -2 | $P(42)$ | 1254.8133 | 11 |
| $P(139)$ | 1245.4511 | -14 | $P(89)$ | 1250.4230 | -5 | $P(41)$ | 1254.9029 | 6 |
| $P(138)$ | 1245.5539 | -1 | $P(88)$ | 1250.5197 | -3 | $P(40)$ | 1254.9938 | 17 |
| $P(137)$ | 1245.6538 | -18 | $P(87)$ | 1250.6146 | -18 | $P(39)$ | 1255.0824 | 5 |
| $P(136)$ | 1245.7571 | 2 | $P(86)$ | 1250.7119 | -8 | $P(38)$ | 1255.1724 | 9 |
| $P(135)$ | 1245.8595 | 12 | $P(85)$ | 1250.8084 | -3 | $P(37)$ | 1255.2619 | 9 |
| $P(134)$ | 1245.9588 | -9 | $P(84)$ | 1250.9041 | -6 | $P(36)$ | 1255.3495 | -8 |
| $P(132)$ | 1246.1625 | 3 | $P(83)$ | 1251.0000 | -6 | $P(35)$ | 1255.4401 | 5 |
| $P(131)$ | 1246.2639 | 7 | $P(82)$ | 1251.0954 | -9 | $P(34)$ | 1255.5297 | 10 |
| $P(130)$ | 1246.3645 | 1 | $P(81)$ | 1251.1900 | -19 | $P(33)$ | 1255.6180 | 4 |
| $P(129)$ | 1246.4661 | 8 | $P(80)$ | 1251.2861 | -12 | $P(32)$ | 1255.7068 | 4 |
| $P(128)$ | 1246.5682 | 20 | $P(79)$ | 1251.3782 | -45 | $P(31)$ | 1255.7953 | 2 |
| $P(127)$ | 1246.6675 | 5 | $P(78)$ | 1251.4765 | -13 | $P(30)$ | 1255.8839 | 3 |
| $P(126)$ | 1246.7681 | 3 | $P(77)$ | 1251.5712 | -17 | $P(29)$ | 1255.9785 | 66 |
| $P(125)$ | 1246.8691 | 6 | $P(76)$ | 1251.6663 | -15 | $P(28)$ | 1256.0604 | 2 |
| $P(124)$ | 1246.9693 | 2 | $P(75)$ | 1251.7599 | -26 | $P(27)$ | 1256.1479 | -4 |
| $P(122)$ | 1247.1692 | -8 | $P(74)$ | 1251.8548 | -24 | $P(26)$ | 1256.2365 | 2 |
| $P(121)$ | 1247.2708 | 5 | $P(73)$ | 1251.9529 | 12 | $P(25)$ | 1256.3248 | 7 |
| $P(120)$ | 1247.3707 | 2 | $P(72)$ | 1252.0471 | 11 | $P(24)$ | 1256.4110 | -8 |
| $P(119)$ | 1247.4702 | -5 | $P(71)$ | 1252.1422 | 19 | $P(23)$ | 1256.4993 | -1 |
| $P(118)$ | 1247.5708 | 1 | $P(70)$ | 1252.2388 | 45 | $P(22)$ | 1256.5858 | -10 |
| $P(117)$ | 1247.6710 | 3 | $P(69)$ | 1252.3308 | 25 | $P(21)$ | 1256.6706 | -35 |
| $P(116)$ | 1247.7707 | 2 | $P(68)$ | 1252.4228 | 7 | $P(20)$ | 1256.7611 | -1 |
| $P(115)$ | 1247.8707 | 4 | $P(67)$ | 1252.5177 | 20 | $P(19)$ | 1256.8483 | 1 |
| $P(114)$ | 1247.9705 | 6 | $P(66)$ | 1252.6103 | 10 | $P(18)$ | 1256.9351 | 0 |
| $P(113)$ | 1248.0706 | 11 | $P(65)$ | 1252.7040 | 13 | $P(17)$ | 1257.0212 | -6 |
| $P(112)$ | 1248.1694 | 5 | $P(64)$ | 1252.7963 | 4 | $P(16)$ | 1257.1077 | -7 |
| $P(111)$ | 1248.2683 | 1 | $P(63)$ | 1252.8893 | 2 | $P(15)$ | 1257.1935 | -14 |
| $P(110)$ | 1248.3675 | 1 | $P(62)$ | 1252.9827 | 7 | $P(14)$ | 1257.2815 | 3 |
| $P(109)$ | 1248.4676 | 10 | $P(61)$ | 1253.0757 | 8 | $P(13)$ | 1257.3661 | -13 |
| $P(108)$ | 1248.5659 | 4 | $P(60)$ | 1253.1677 | 1 | $P(12)$ | 1257.4531 | -3 |
| $P(107)$ | 1248.6640 | -4 | $P(59)$ | 1253.2602 | 2 | $P(11)$ | 1257.5348 | -45 |
| $P(106)$ | 1248.7631 | 0 | $P(58)$ | 1253.3528 | 3 | $P(10)$ | 1257.6248 | -3 |
| $P(105)$ | 1248.8620 | 2 | $P(57)$ | 1253.4454 | 6 | $P(9)$ | 1257.7121 | 13 |
| $P(104)$ | 1248.9600 | -4 | $P(56)$ | 1253.5368 | -2 | $P(8)$ | 1257.7955 | -7 |
| $P(103)$ | 1249.0583 | -5 | $P(55)$ | 1253.6289 | 0 | $P(7)$ | 1257.8792 | -24 |
| $P(102)$ | 1249.1552 | -19 | $P(54)$ | 1253.7215 | 7 | $P(6)$ | 1257.9625 | -43 |
| $P(101)$ | 1249.2539 | -13 | $P(53)$ | 1253.8108 | -17 | $P(5)$ | 1258.0506 | -14 |
| $P(100)$ | 1249.3526 | -6 | $P(52)$ | 1253.9049 | 9 | $P(4)$ | 1258.1367 | -2 |
| $P(99)$ | 1249.4490 | -21 | $P(51)$ | 1253.9938 | -17 | $P(3)$ | 1258.2158 | -59 |
| $P(98)$ | 1249.5476 | -14 | $P(50)$ | 1254.0868 | 0 | $P(2)$ | 1258.3060 | -4 |
| $P(97)$ | 1249.6489 | 23 | $P(49)$ | 1254.1764 | -16 | $R(0)$ | 1258.5586 | -10 |
| $P(96)$ | 1249.7448 | 6 | $P(48)$ | 1254.2649 | -41 | $R(1)$ | 1258.6400 | -37 |
| $P(95)$ | 1249.8409 | -7 | $P(47)$ | 1254.3594 | -5 | $R(2)$ | 1258.7253 | -24 |
| $P(94)$ | 1249.9385 | -4 | $P(46)$ | 1254.4502 | -4 | $R(3)$ | 1258.8090 | -26 |
| $P(93)$ | 1250.0344 | -17 | $P(45)$ | 1254.5420 | 8 | $R(4)$ | 1258.8938 | -15 |
| $P(92)$ | 1250.1328 | 4 | $P(44)$ | 1254.6334 | 18 | $R(5)$ | 1258.9819 | 31 |
| $R(6)$ | 1259.0618 | -5 | $R(37)$ | 1261.5807 | 1 | $R(68)$ | 1263.9685 | 43 |
| $R(7)$ | 1259.1431 | -25 | $R(38)$ | 1261.6610 | 14 | $R(69)$ | 1264.0432 | 44 |
| $R(8)$ | 1259.2277 | -10 | $R(39)$ | 1261.7395 | 10 | $R(70)$ | 1264.1157 | 24 |
| $R(9)$ | 1259.3120 | 2 | $R(40)$ | 1261.8182 | 10 | $R(71)$ | 1264.1925 | 48 |
| $R(10)$ | 1259.3933 | -15 | $R(41)$ | 1261.8965 | 7 | $R(72)$ | 1264.2650 | 31 |
| $R(11)$ | 1259.4753 | -22 | $R(42)$ | 1261.9743 | 1 | $R(74)$ | 1264.4080 | -18 |
| $R(12)$ | 1259.5590 | -11 | $R(43)$ | 1262.0541 | 16 | $R(75)$ | 1264.4821 | -15 |
| $R(13)$ | 1259.6425 | -1 | $R(44)$ | 1262.1311 | 4 | $R(76)$ | 1264.5557 | -15 |
| $R(14)$ | 1259.7250 | 1 | $R(45)$ | 1262.2068 | -20 | $R(77)$ | 1264.6317 | 10 |
| $R(15)$ | 1259.8069 | -3 | $R(46)$ | 1262.2844 | -22 | $R(78)$ | 1264.7058 | 17 |
| $R(16)$ | 1259.8887 | -5 | $R(47)$ | 1262.3642 | -2 | $R(79)$ | 1264.7775 | 2 |
| $R(17)$ | 1259.9706 | -5 | $R(48)$ | 1262.4413 | -6 | $R(80)$ | 1264.8526 | 22 |
| $R(18)$ | 1260.0534 | 5 | $R(49)$ | 1262.5188 | -7 | $R(81)$ | 1264.9215 | -19 |
| $R(19)$ | 1260.1345 | -1 | $R(50)$ | 1262.5973 | 6 | $R(82)$ | 1264.9950 | -12 |
| $R(20)$ | 1260.2168 | 7 | $R(51)$ | 1262.6726 | -13 | $R(83)$ | 1265.0673 | -16 |
| $R(21)$ | 1260.2981 | 6 | $R(52)$ | 1262.7506 | -4 | $R(84)$ | 1265.1411 | -3 |
| $R(22)$ | 1260.3788 | 1 | $R(53)$ | 1262.8277 | -2 | $R(85)$ | 1265.2183 | 46 |
| $R(23)$ | 1260.4605 | 7 | $R(54)$ | 1262.9042 | -4 | $R(86)$ | 1265.2838 | -22 |

TABLE II. (Continued.)

| Line | OBS | $\begin{gathered} \mathrm{OMC}^{\mathrm{a}} \\ \times 10^{4} \end{gathered}$ | Line | OBS | $\begin{aligned} & \mathrm{OMC} \\ & \times 10^{4} \end{aligned}$ | Line | OBS | $\begin{gathered} \mathrm{OMC} \\ \times 10^{4} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $R(24)$ | 1260.5413 | 5 | $R(55)$ | 1262.9810 | -2 | $R(87)$ | 1265.3573 | -9 |
| $R(25)$ | 1260.6240 | 24 | $R(56)$ | 1263.0565 | -12 | $R(88)$ | 1265.4311 | 10 |
| $R(26)$ | 1260.7024 | 1 | $R(57)$ | 1263.1349 | 9 | $R(89)$ | 1265.5064 | 44 |
| $R(27)$ | 1260.7838 | 10 | $R(58)$ | 1263.2101 | -1 | $R(90)$ | 1265.5734 | -3 |
| $R(28)$ | 1260.8644 | 11 | $R(59)$ | 1263.2857 | -5 | $R(91)$ | 1265.6459 | 7 |
| $R(29)$ | 1260.9450 | 15 | $R(60)$ | 1263.3637 | 16 | $R(92)$ | 1265.7157 | -10 |
| $R(30)$ | 1261.0247 | 10 | $R(61)$ | 1263.4382 | 3 | $R(93)$ | 1265.7862 | -18 |
| $R(31)$ | 1261.1037 | 0 | $R(62)$ | 1263.5143 | 9 | $R(94)$ | 1265.8591 | -2 |
| $R(32)$ | 1261.1837 | 3 | $R(63)$ | 1263.5893 | 4 | $R(95)$ | 1265.9321 | 18 |
| $R(33)$ | 1261.2644 | 12 | $R(64)$ | 1263.6646 | 3 | $R(96)$ | 1266.0000 | -13 |
| $R(34)$ | 1261.3437 | 9 | $R(65)$ | 1263.7410 | 16 | $R(97)$ | 1266.0745 | 24 |
| $R(35)$ | 1261.4228 | 6 | $R(66)$ | 1263.8161 | 15 | $R(98)$ | 1266.1414 | -14 |
| $R(36)$ | 1261.5028 | 13 | $R(67)$ | 1263.8925 | 31 | $R(99)$ | 1266.2132 | -1 |

${ }^{a}$ Observed minus calculated line positions.
for these lines. However, many of the weaker lines and the blended features were determined only to a precision of about $\pm 0.005 \mathrm{~cm}^{-1}$.

An interactive color Loomis-Wood computer program was used to pick out the branches and helped to assign the spectra. The CO molecule was present in the cell as a side product of the thermolysis. The measured spectral lines of the $\nu_{4}$ vibrational mode were calibrated with the CO lines using line positions taken from the literature. ${ }^{9}$ For the $\nu_{5}$ vibrational mode, we used the water absorption for calibration. ${ }^{10}$

## The $\nu_{4}$ fundamental mode

Figure 1 shows the high-resolution spectrum of the CNO antisymmetric stretching mode $\left(\nu_{4} \sigma_{u}\right)$ of ONCCNO. The HNCO molecule is present on the high wave number side of the overview spectrum [Fig. 1(a)]. Figure 1(b) is an expanded portion of the $R$ branch of the $\nu_{4}$ fundamental. The alternation in intensities is caused by the nuclear spin statistics of two equivalent $(I=1)$ nitrogen atoms. The symmetrically located nitrogen atoms cause an intensity alternation of 2:1 with the even $J^{\prime \prime}$ values being stronger.

The main problem in the analysis was the presence of a large number of hot bands which causes a line density as high as 100 lines $/ \mathrm{cm}^{-1}$. It was impossible to locate the band origin because of the overlapping hot bands. The final $J$ assignments were achieved with the help of a small local perturbation at $J^{\prime}=97$ in the upper level. Figure 2 shows the perturbed $R(96)$ and $P(98)$ lines. The $R(96)$ line was shifted by $0.005 \mathrm{~cm}^{-1}$ to higher wave numbers while the $P(98)$ line was shifted by the same amount. There are several other small local perturbations in this band which also confirm the assignments. In total, 230 rotational lines were assigned for the $\nu_{4}$ band and they are listed in Table I.

## The $\nu_{5}$ fundamental mode

The high-resolution spectrum of the CNO symmetric stretching mode ( $\nu_{5} \sigma_{u}$ ) is shown in Fig. 3(a). An expanded portion of the $P$ branch is shown in Fig. 3(b), again demonstrating the $2: 1$ intensity alternation due to nitrogen nuclear spin statistics. The measurement and assignment of $\nu_{5}$ pro-
ceeded in the same way as for $\nu_{4}$. The $\nu_{5}$ mode also had a high line density but no perturbations were present to help the assignment. The rotational assignment of $\nu_{5}$ was made by changing the relative assignment of the $P$ and $R$ branches until the ground state combination differences matched those of $\nu_{4}$. About 200 lines were assigned and the complete line list is shown in Table II.

## The rotational constants

The line positions (Tables I and II) of the two fundamental bands were fitted together in a global least-squares fit. The energy level expression

$$
F(J)=\nu_{0}+B J(J+1)-D[J(J+1)]^{2}+H[J(J+1)]^{3}
$$

was used in the fit and the resulting rotational constants are listed in Table III. The ground state $B_{0}$ constant, $0.04220210 \mathrm{~cm}^{-1}(1.265187 \mathrm{GHz})$, is in good agreement with the $a b$ initio prediction ${ }^{8}$ of $1.24 \sim 1.25 \mathrm{GHz}$. The centrifugal distortion constants of the excited $\nu_{4}$ vibrational level may be perturbed by interaction with other modes.

## DISCUSSION AND CONCLUSION

The semistable cyanogen di-N-oxide (ONCCNO) molecule can be generated in situ with a good yield in the gas phase from the precursor molecule dichloroglyoxime by thermolysis. The high-resolution spectra of the ONCCNO molecule have been obtained for the two strongest vibrational modes $\nu_{4}$ and $\nu_{5}$. The rotational analysis of these two modes and the intensity alternation of the lines indicates that the ONCCNO molecule is not bent.

TABLE III. The rotational constants for ONCCNO (in $\mathrm{cm}^{-1}$ ). One standard deviation is in parentheses.

| Level | $\nu_{0}$ | $B$ | $10^{9} \mathrm{D}$ | $10^{14} \mathrm{H}$ |
| :--- | :---: | :---: | :---: | :---: |
| Ground | 0.0 | $0.04220210(96)$ | $0.877(70)$ | $\ldots$ |
| $\nu_{4}$ | $1258.47530(11)$ | $0.04213364(98)$ | $1.071(74)$ | $1.196(38)$ |
| $\nu_{5}$ | $2246.04055(23)$ | $0.04210939(97)$ | $1.978(73)$ | $3.503(96)$ |

TABLE IV. Estimated geometry of ONCCNO. ${ }^{\text {a }}$

| Parameter | Structure I | Structure II |
| :--- | :---: | :---: |
| $r_{\mathrm{NO}} / \AA$ | $1.1994^{\mathrm{b}}$ | $1.1923^{\mathrm{c}}$ |
| $r_{\mathrm{CN}} / \AA$ | $1.1679^{\mathrm{b}}$ | $1.1730^{\mathrm{c}}$ |
| $r_{\mathrm{CC}} / \AA$ | $1.3337^{\mathrm{d}}$ | $1.3329^{\mathrm{d}}$ |
| $\angle \mathrm{ONC}$ | $180.0^{\circ}$ | $180.0^{\circ}$ |
| $\angle \mathrm{NCC}$ | $180.0^{\circ}$ | $180.0^{\circ}$ |

${ }^{\text {a }}$ Linear framework assumed.
${ }^{\mathrm{b}} r_{S}$ bond lengths from HCNO, Ref. 11.
${ }^{\mathrm{c}}$ See text for details.
${ }^{\mathrm{d}}$ Bond length calculated using experimental rotational constant.

Indications from the $a b$ initio calculations, and comparisons with similar molecules, suggest that the important CC bond is quite short. ${ }^{8}$ The single $B$ value obtained in this work is insufficient to provide this parameter without some assumptions about the CN and NO bond lengths. One approach to this problem is to use the known CN and NO bond lengths in the parent HCNO molecule, ${ }^{11}$ and assuming a linear structure, extract the $C C$ value from the rotational constant $\left(B_{0}\right)$ determined in this work. Table IV, structure I, shows the structure of ONCCNO, predicted by this method. This, of course, assumes that the CN and NO bond lengths are transferable from HCNO .

A possible improvement on this method is to see how a computational method $\left[\mathrm{MP} 3(f \mathrm{full}) / 6-31 \mathrm{G}^{*}\right]^{8}$ performs for both HCNO and ONCCNO, and then "correct" the CN and NO values and apply them to ONCCNO. With these "refined" CN and NO bond lengths, the mutual effect of back-to-back CNO groups should be taken into account, and, using the experimental rotational constant, a value for the CC bond length is obtained. This assumes that the differences between the CN and NO bond lengths in HCNO and ONCCNO are predicted correctly by the MP3 method. Table IV, structure II, shows the revised values for ONCCNO. This structure which takes into account the interaction of the two CNO groups, indicates that the NO bond length has slightly decreased, while the CN bond length has slightly increased, in accord with expectations based on a more delocalized framework.

Both the directly transferred CN and NO lengths, and the corrected values, lead to the result that the CC bond length is about $1.333 \AA$, approximately the same as that in ethylene, $\mathrm{H}_{2} \mathrm{CCH}_{2}$, a molecule of known double bond character. Table V shows CC bond lengths for a range of molecules. An important observation is that the experimental CC bond length in cyanogen, NC-CN is $1.389 \AA,{ }^{12}$ already much shorter than a typical CC single bond and we anticipate that this will decrease upon adding terminal oxygen atoms (compare $\mathrm{H}_{3} \mathrm{CCN}$ and $\mathrm{H}_{3} \mathrm{CCNO}$ in Table V ). We note that MP3(full)/6-31G* calculations for NCCN predict CC to be $1.393 \AA$, with similar calculations for ONCCNO predicting CC to be $1.361 \AA .{ }^{8}$ The inference from this is that the CC length in ONCCNO is at least $0.03 \AA$ less than that in NCCN.

Clearly, experimental and theoretical studies on this molecule have a long way to go before a full characterization is

TABLE V. The $\mathrm{C}-\mathrm{C}$ bond lengths $\left(r_{0}\right)$ in simple organic molecules.

| Molecule | Method | $r_{\mathrm{C}-\mathrm{C}}(\AA)$ |
| :--- | :---: | :--- |
| $\mathrm{H}_{3} \mathrm{C}-\mathrm{CH}_{3}$ | IR | $1.536^{\mathrm{a}}$ |
| $\mathrm{H}_{3} \mathrm{C}-\mathrm{C} \equiv \mathrm{N}$ | MW | $1.4582^{\mathrm{b}}$ |
| $\mathrm{H}_{3} \mathrm{C}-\mathrm{C} \equiv \mathrm{NO}$ | MW | $1.442^{\mathrm{c}}$ |
| $\mathrm{N} \equiv \mathrm{C}-\mathrm{C} \equiv \mathrm{N}$ | IR | $1.389^{\mathrm{d}}$ |
| $\mathrm{HC} \equiv \mathrm{C}-\mathrm{C} \equiv \mathrm{N}$ | MW | $1.382^{\mathrm{e}}$ |
| $\mathrm{H}_{2} \mathrm{C}=\mathrm{CH}$ | IR | $1.337^{\mathrm{f}}$ |
| $\mathrm{ONC}-\mathrm{CNO}$ | IR | $1.3329^{\mathrm{g}}$ |
| $\mathrm{HC} \equiv \mathrm{CH}$ | IR | $1.20862^{\mathrm{h}}$ |

${ }^{\text {a }}$ Reference 13.
frefer 16
${ }^{\circ}$ Reference 15 . ${ }^{\mathrm{g}}$ This work.
${ }^{\mathrm{d}}$ Reference $12 . \quad{ }^{\mathrm{h}}$ Reference 18.
achieved. For example, the possibility of quasilinear behavior cannot be ruled out at this stage of the analysis. We plan additional experiments to record spectra of the much weaker bending vibrational modes as well as combination bands. This work will help assign the numerous hot bands that we have measured and will provide information on the missing modes of gerade symmetry.

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