Fourier Transform Infrared Emission Spectroscopy of the $C^4\Delta - X^4\Phi$ System of ZrCl

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Received July 18, 1997; in revised form September 12, 1997

The high-resolution spectrum of ZrCl has been investigated in emission in the region $3000-10\,000~\rm cm^{-1}$ using a Fourier transform spectrometer. The bands were excited in a microwave discharge through a flowing mixture of ZrCl₄ and helium. New ZrCl bands observed in the interval $3600-4400~\rm cm^{-1}$ have been assigned to a new $C^4\Delta-X^4\Phi$ electronic transition. Four bands with *R* heads at 4305.7, 4208.0, 4064.7, and 3897.4 cm⁻¹ have been assigned as the 0–0 bands of the $^4\Delta_{1/2}-^4\Phi_{3/2}$, $^4\Delta_{3/2}-^4\Phi_{5/2}$, $^4\Delta_{5/2}-^4\Phi_{7/2}$, and $^4\Delta_{7/2}-^4\Phi_{9/2}$ subbands. A rotational analysis of the 0–0, 1–1, and 2–2 vibrational bands of the four subbands has been obtained and molecular constants have been extracted. Although there is no proof that the lower $^4\Phi$ state is the ground state of ZrCl, we have labeled this transition as $C^4\Delta-X^4\Phi$, consistent with the corresponding near-infrared transition of TiCl (R. S. Ram and P. F. Bernath, *J. Mol. Spectrosc.*, in press). © 1997 Academic Press

INTRODUCTION

The emission spectrum of ZrCl was initially observed by Carroll and Daly (1) from the radiofrequency excitation of ZrCl₄ vapor obtained by heating an anhydrous sample of ZrCl₄ powder. Many complex bands were observed in the region 280-420 nm which were tentatively classified into three electronic transitions named as system A (280-290 nm), system B (360-380 nm), and system C (400-415 nm). A tentative electronic assignment of ${}^{4}\Pi - {}^{4}\Sigma^{-}$ was suggested for system C on the basis of a vibrational analysis, although no rotational analysis was presented. The emission spectrum of ZrCl has also been recently studied by Jordan et al. (2), who excited the molecule in a low-temperature corona-excited supersonic jet discharge. They classified the bands in the region 400-415 nm into four groups of double-headed bands assigned as the ${}^4\Pi_{-1/2} - {}^4\Sigma^-$, ${}^4\Pi_{1/2} - {}^4\Sigma^-$, ${}^4\Pi_{3/2} - {}^4\Sigma^-$, and ${}^{4}\Pi_{5/2} - {}^{4}\Sigma^{-}$ subbands of a Hund's case (a) ${}^{4}\Pi$ -Hund's case (b) ${}^{4}\Sigma^{-}$ transition. Again no rotational analysis of the bands was provided to support their assignments. Another transition of ZrCl has been observed in the near infrared by Phillips et al. (3), who reported a rotational analysis of numerous bands and concluded that these bands were probably due to a ${}^{2}\Pi - {}^{2}\Pi$ transition (although $^{2}\Delta - ^{2}\Delta$ and $^{2}\Phi - ^{2}\Phi$ assignments were not excluded). A definite assignment of the ground electronic state of ZrCl is still lacking. Sivaji and Rao have found violet and nearultraviolet bands of ZrBr (4) and ZrI (5), similar to the Carroll and Daly systems of ZrCl (1).

There are no theoretical calculations for any of the Zr-containing halides to assist in the assignment of the observed spectra. However, an ab initio calculation for TiF has recently been carried out by Harrison (6) and some work is in progress at the University of Utah (7). The work of the Harrison group predicts a ${}^4\Phi$ ground state for TiF and also provides spectroscopic properties for several low-lying electronic states (6). Our recent observation of the $G^4\Phi - X^4\Phi$ transition of TiF near 15 000 cm⁻¹ (8) is consistent with a $^4\Phi$ ground state. Recently we have also investigated the electronic spectra of TiCl in the region $3000-12500 \text{ cm}^{-1}(9)$ and have found three new electronic transitions which have been named $C^4\Delta$ - $X^{4}\Phi$, $G^{4}\Phi - X^{4}\Phi$, and $G^{4}\Phi - C^{4}\Delta$. We have based our letter notation for the different electronic states of TiCl on the theoretical predictions for TiF (6) and the data available for the isovalent TiH (10-12).

In the present work we report on the results of a search for the electronic transitions of ZrCl in the region $3000-10\,000\,\mathrm{cm^{-1}}$. We have observed a $^4\Delta-^4\Phi$ electronic transition of ZrCl in the region $3700-4400\,\mathrm{cm^{-1}}$ which is analogous to the $C^4\Delta-X^4\Phi$ transition of TiCl (9). A rotational analysis of a number of vibrational bands of different subbands has been obtained and the spectroscopic data for all four spin components of the $X^4\Phi$ and $C^4\Delta$ states have been extracted. Several weaker bands and some complex bands remain to be assigned. Although we do not have any proof for our assignment of the lower state as the ground state, we have decided to maintain

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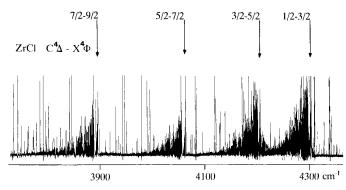


FIG. 1. A compressed portion of the $C^4\Delta - X^4\Phi$ transition of ZrCl with the *R* heads marked for the 0–0 vibrational bands of each subband.

the $C^4\Delta - X^4\Phi$ notation for this transition of ZrCl to be consistent with the recent results on the isovalent TiCl molecule.

EXPERIMENTAL DETAILS

The ZrCl molecules were excited in an electrodeless microwave discharge through a flowing mixture of 3 Torr of He and a trace of ZrCl₄. The discharge tube was made of quartz and had an outer diameter of 12 mm. A small glass bulb was attached to the discharge tube with a $\frac{1}{2}$ -in. Cajon Ultratorr fitting. A small quantity of solid ZrCl₄ was placed in the bulb, which was constantly heated with a heat gun to maintain a blue-white discharge. The He gas flowed over the heated ZrCl₄ sample and carried the vapor into the discharge region. The emission from the discharge tube passed directly through the 8-mm entrance aperture of the 1-m Fourier transform spectrometer of the National Solar Observatory at Kitt Peak. The spectra in the interval 1800–9000 cm⁻¹ were initially recorded in the first-order alias using liquid nitrogen-cooled InSb detectors and Si filters. A total of 10 scans were coadded in about 60 min of integration at a resolution of 0.02 cm⁻¹. The region 3000-5000 cm⁻¹ was also recorded in the third-order alias at a resolution of 0.01 cm⁻¹. We have used this spectrum in the analysis of the $C^4\Delta$ - $X^4\Phi$ transition.

The spectral line positions were determined using a data reduction program called PC-DECOMP developed by J. Brault. The peak positions were determined by fitting a Voigt lineshape function to each line. The spectra were calibrated using the wavenumbers of the vibration–rotation lines of the $1{\text -}0$ band of HCl (13) which also appeared in emission in the same spectrum. The molecular lines of ZrCl had a typical width of $0.015~\text{cm}^{-1}$ and appeared with a maximum signal-to-noise ratio of $15{:}1$ so that the best line positions are expected to be accurate to about $\pm 0.002~\text{cm}^{-1}$.

OBSERVATION AND ANALYSIS

The present spectrum contains a number of new ZrCl bands in the region 3000–10 000 cm⁻¹. Several of the new bands observed in the interval 6000–10 000 cm⁻¹ belong to the two subbands $({}^{2}\Pi_{1/2} - {}^{2}\Pi_{1/2})$ and ${}^{2}\Pi_{3/2} - {}^{2}\Pi_{3/2})$ previously studied by Phillips *et al.* (3). In the region $3000-4500 \text{ cm}^{-1}$, however, there are four prominent groups of $|\Delta\Omega| = 1$ bands with strong R heads at 4305.7, 4208.0, 4064.7, and 3897.4 cm^{-1} which have been assigned as the 0-0 bands of the ${}^4\Delta_{1/2} - {}^4\Phi_{3/2}$, ${}^4\Delta_{3/2} - {}^4\Phi_{5/2}$, ${}^4\Delta_{5/2} - {}^4\Phi_{7/2}$, and ${}^4\Delta_{7/2} ^4\Phi_{9/2}$ subbands, respectively, of the $C^4\Delta - X^4\Phi$ transition. This assignment is consistent with our recent work on the analogous transition of TiCl in the region 3100–3400 cm⁻¹. A portion of the compressed spectrum of the $C^4\Delta - X^4\Phi$ transition of ZrCl is presented in Fig. 1. Each of the 0-0bands in Fig. 1 is followed to lower wavenumbers by the 1-1, 2-2, and 3-3 bands with decreasing intensity. Offdiagonal bands were not identified in our spectra because of their very weak intensity; the vibrational intervals remain to be determined.

The rotational structure of each band consists of P, Q, and R branches. The Q branch is the strongest and the P branch is stronger than the R branch, consistent with a $\Delta\Omega$ = -1 transition. The lines of the ${}^4\Delta_{1/2} - {}^4\Phi_{3/2}$ subband are doubled by the presence of appreciable Ω -doubling in the excited $^4\Delta_{1/2}$ state. The lines of the $^4\Delta_{3/2}$ $^4\Phi_{5/2}$, $^4\Delta_{5/2}$ $^4\Phi_{7/2},$ and $^4\Delta_{7/2}-^4\Phi_{9/2}$ subbands, however, do not show any Ω -doubling. Zr has five isotopes, 90 Zr (51.5%), 91 Z (11.2%), 92 Zr (17.2%), 94 Zr (17.4%), and 96 Zr (2.7%), but we have determined the molecular constants only for the most abundant 90Zr35Cl isotopomer. Some weak lines of the minor isotopomer 90Zr37Cl have also been observed for the strong bands but the data were not sufficient for an independent rotational analysis. A band with weak doubled R heads at 4309.3 and 4309.6 cm⁻¹ was also found near the 0-0 band of the ${}^4\Delta_{1/2}$ – ${}^4\Phi_{3/2}$ subband, but it probably belongs to another transition. Similarly, some rotational structure without any prominent head has been observed on the high wavenumber

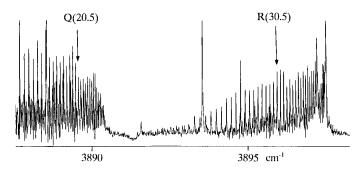


FIG. 2. An expanded portion of the 0-0 band of the $C^4\Delta_{7/2}-X^4\Phi_{9/2}$ subband of ZrCl near the R head.

TABLE 1 Observed Line Positions (in cm $^{-1}$) for the $C^4\Delta$ – $X^4\Phi$ Transition of ZrCl

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	70.5	4303.343	-3	4270.236	-1	4304.105	-12	4270.993	1	4286.674	-2	4287.436	-3	70.5

Note. O-C are observed minus calculated line positions in the units of 10^{-3} cm⁻¹ and asterisks mark perturbed lines.

TABLE 1—Continued

						-Continue						
\mathbf{I} \mathbf{R}_{ee}	О-С	P _{ee}	О-С	\mathbf{R}_{ff}	О-С	\mathbf{P}_{ff}	О-С	$\mathbf{Q}_{ ext{ef}}$	О-С	Q_{fe}	О-С	J
71.5 4303.219 72.5 4303.084 73.5 4302.959 74.5 4302.812 75.5 4302.521 77.5 4302.215 79.5 4302.215 79.5 4301.891 81.5 4301.716 82.5 4301.541 83.5 4301.174 85.5 4300.987 86.5 4300.987 86.5 4300.997 87.5 4300.997 87.5 4300.997 87.5 4300.179 90.5 4299.965 91.5 4299.753 92.5 94.5	-1 -5 4 -4 -2 -2 0 2 -3 7 3 4 2 2 4 8 -1 -8 -3 -6 -3	4269.647 4269.053 4268.468 4267.856 4267.254 4266.029 4265.411 4264.786 4264.163 4263.531 4262.899 4262.257 4261.613 4259.655 4258.992 4258.325 4257.656 4256.988 4256.300 4255.619 4254.935	-2 -3 8 -3 1 -5 -1 -1 -3 1 0 3 2 2 2 4 2 0 -2 -1 4 -8 -9 -10	4303.994 4303.616 4303.470 4303.343 4303.192 4303.043 4302.729 4302.564 4302.394 4302.215 4302.048 4301.861 4301.673 4301.487 4301.290 4301.083 4300.881 4300.672 4300.052 4300.227 4300.005 4299.782	-5 -1 -10 3 -2 0 -2 0 1 0 -6 6 3 2 8 8 2 6 7 0 -7 -7 -5	4270.404 4269.821 4269.233 4268.639 4268.044 4267.442 4266.839 4266.225 4265.612 4264.984 4263.743 4263.743 4263.181 4260.530 4259.220 4258.557 4257.887 4257.887 4257.211 4256.534	-7 -6 -6 -4 -4 0 -9 -5 4 1 4 5 2 1 7 6 6 3 -1 -3	4286.317 4285.953 4285.592 4285.221 4284.846 4284.465 4283.696 4283.303 4282.907 4282.506 4282.099 4281.692 4281.278 4280.861 4280.439 4280.439	-2 -4 0 -1 -1 -3 0 -1 -2 -1 -1 2 1 4 5 7	4287.086 4286.732 4286.378 4286.011 4285.644 4285.278 4284.520 4284.133 4283.743 4283.350 4282.952 4282.550 4282.143 4281.313 4281.313 4280.890 4280.467 4280.038	-3 -4 0 -4 -5 1 -4 -1 -3 -2 -1 0 2 3 5 3 0 4 5	71.5 72.5 73.5 74.5 75.5 76.5 77.5 80.5 81.5 82.5 83.5 84.5 85.5 86.5 89.5 90.5 91.5 92.5 93.5 94.5
				C ⁴ /	$\Delta_{1/2}$ - X	C ⁴ Φ _{3/2} 1-	1					
6.5 7.5 8.5 9.5 4282.333 10.5 4282.669 11.5 4282.669 12.5 4283.303 16.5 17.5 4283.303 16.5 17.5 4283.822 20.5 4284.361 21.5 4284.059 22.5 4284.361 25.5 26.5 4284.361 25.5 26.5 4284.618 28.5 29.5 4284.618 28.5 29.5 4284.888 32.5 4284.618 28.5 29.5 4284.898 32.5 4284.898 32.5 4284.960 31.5 4285.068 35.5 4285.093 36.5 4285.093 36.5 4285.190 39.5 40.5 41.5 42.5 43.5 44.5	-5 -3 -8 -2 1 4 11 14 0 2 1 5 -6 -5 -2 -5 -7 -11 5 -4 8 -3 12 -3 7 2 0	4278.530 4278.251 4277.975 4277.669 4277.686 4276.786 4276.157 4275.827 4275.827 4275.827 4275.827 4273.803 4274.503 4273.441 4273.081 4273.441 4273.081 4272.705 4272.333 4271.579 4270.404 4269.993 4269.993 4269.983 4266.8329 4267.897 4267.897 4266.8329 4267.897 4266.8329 4267.897 4266.8329 4267.897 4266.8329 4267.897 4266.585 4266.135 4266.135 4264.764 4264.293	1 0 6 -11 0 1 4 -2 0 9 0 1 -3 2 9 2 -3 0 -8 -7 1 -5 1 7 -1 -2 -4 5 2 -2 1 6 -3 6 3	4282.054 4282.228 4282.416 4282.607 4282.780 4282.952 4283.099 4283.426 4283.581 4283.851 4284.350 4284.231 4284.350 4284.461 4284.846 4284.943 4285.276 4285.276 4285.343 4285.398 4285.398 4285.398 4285.398	6 -11 -10 1 -1 2 -16 -2 5 -5 -4 -10 -6 -3 -3 -3 -4 -2 -10 2 -1 2 -11 3 9 7 0 3 5 4	4276.883 4276.575 4276.266 4275.960 4275.641 4275.317 4274.988 4274.650 4274.312 4273.263 4272.542 4272.160 4271.793 4271.024 4270.637 4270.236 4269.843 4269.436 4269.436 4269.436 4269.436 4269.436 4269.436 4269.436 4269.436 4269.436 4269.55 4266.8185 4267.324 4266.883 4266.450 4266.450 4265.551 4265.097 4264.639	5 -1 -3 4 3 3 2 -2 -1 3 3 -9 0 -10 -2 2 -3 5 4 4 11 2 0 -2 -7 1 2 0 1 4	4279.947 4279.897 4279.690 4279.623 4279.544 4279.448 4279.259 4279.259 4279.259 4278.813 4278.685 4278.555 4278.813 4278.685 4278.279 4277.329 4277.329 4277.329 4277.329 4277.329 4277.501 4276.996 4276.397 4275.799 4275.799 4275.799 4275.799 4275.799 4275.799 4275.799 4275.799 4275.799 4275.799 4275.799 4275.340 4275.340 4275.340 4275.340 4275.340 4275.340 4275.340 4275.340 4275.340 4275.340	0 5 12 8 -4 6 9 0 -1 3 5 -2 2 -1 -2 1 -6 3 0 -1 -2 0 -3 -2 0 -3 -2 -1 -1 1 -2 0 5 2	4280.014 4279.971 4279.922 4279.731 4279.657 4279.401 4279.401 4279.206 4279.206 4279.206 4279.206 4278.746 4278.875 4278.875 4278.876 4278.876 4278.210 4278.251 4277.588 4277.748 4277.588 4277.684 4277.684 4277.684 4276.693 4276.094 4275.886 4276.497 4276.586 4275.672 4275.886	-6 -1 3 -1 3 2 1 0 1 2 2 3 4 5 -2 -1 1 0 2 -3 1 0 2 0 8 0 2 1 -1 1 -1 0 0 13	6.5 7.5 8.5 9.5 10.5 11.5 13.5 14.5 15.5 16.5 17.5 18.5 20.5 22.5 23.5 24.5 22.5 23.5 24.5 23.5 30.5 31.5 33.5 34.5 33.5 34.5 35.5 36.5 37.5 38.5 38.5 38.5 38.5 38.5 38.5 38.5 38

TABLE 1—Continued

						–Continue						
J R _{ee}	O-C	\mathbf{P}_{ee}	O-C	\mathbf{R}_{ff}	О-С	$\mathbf{P}_{\mathbf{ff}}$	О-С	\mathbf{Q}_{ef}	О-С	Q_{fe}	О-С	J
45.5 46.5 47.5 48.5 50.5 51.5 52.5 53.5 54.5 55.5 56.5 57.5 58.5 59.5 60.5 61.5 62.5 63.5 64.5 65.5 66.5 77.5 78.5 79.5 80.5 81.5 82.5 83.5 84.5 85.5 86.5 87.5 88.5 89.5		4263.821 4263.344 4262.858 4262.370 4261.879 4261.382 4260.877 4260.370 4259.862 4259.342 4258.298 4257.764 4257.228* 4256.138* 4255.587* 4255.587* 4255.426* 4253.879* 4244.65* 4244.086* 4244.086* 4244.086* 4244.086* 4244.368 4243.036 4242.388 4241.069 4240.411 4239.732 4239.060 4238.380 4237.695	-11 -17 -20 -33 -36 -57 42 35 33 24 19 16			4264.163 4263.699 4263.228 4262.745 4262.257 4261.770 4261.275 4260.268 4259.759 4259.248 4258.209 4257.675* 4255.503* 4256.052* 4253.813* 4253.235* 4250.446* 4249.855* 4244.0371 4246.797 4246.797 4246.177 4245.552 4244.919 4244.283 4243.638 4242.997 4242.351 4241.699 4241.033 4240.371 4239.704 4239.032 4238.354 4237.669 4236.988	33 34 27 22 15	4274.406 4274.162 4273.910 4273.655 4273.397 4273.128 4272.861 4272.582 4272.304 4271.726 4271.432 4271.132 4270.201* 4269.876* 4269.876* 4267.254* 4266.833* 4266.135* 4265.754* 4264.591* 4264.200 4263.395 4264.591* 4265.754* 4266.138* 4267.764* 4268.858*	-17 -24 -33 -44 -63	4274.766 4274.529 4274.284 4274.040 4273.784 4273.531 4273.263 4272.721 4272.442 4272.160 4271.579 4270.3503 4270.0263 4270.0263 4269.3473 4268.9943 4266.6773 4266.6773 4266.6773 4266.3093 4261.505 4264.387 4262.769 4262.769 4262.769 4262.769 4262.769 4262.769 4262.769 4262.769 4262.769 4262.769 4262.769 4262.769 4263.583 4263.583 4263.583 4263.79 4262.769 4260.642 4260.642 4260.642 4260.642 4260.642 4260.206 4259.759 4259.308 4258.852 4258.396	-14 -21 -26 -36 -36 -47 -70 -95 -65 -60 -41 -31 -31 -31 -31 -31 -31 -31 -3	45.5 46.5 47.5 48.5 50.5 51.5 52.5 53.5 55.5 56.5 57.5 60.5 61.5 62.5 64.5 65.5 66.5 67.5 77.5 77.5 77.5 77.5 77
				C ⁴ / ₂	Δ _{1/2} - X	$4\Phi_{3/2}$ 2-2	2					
8.5 9.5 10.5 11.5 12.5 4262.158 13.5 4262.328 14.5 4262.473 15.5 16.5 4262.769 17.5 4262.899 18.5 19.5 4263.160 20.5 4263.290	9 -1 4 -4	4255.814 4255.503 4254.878 4254.549 4253.879 4253.533	-10 -11 1 -2 -6 -10	4262.257 4262.417 4262.573 4262.883 4263.160 4263.290 4263.422	8 3 -1 4 -3 -7 -4	4257.358 4257.073 4256.795 4256.503 4256.207 4255.904 4254.979 4254.658 4254.335 4254.002 4253.674	8 1 7 3 1 -2 0 -1 1 -3 4	4259.278 4259.220 4259.163 4259.100 4259.037 4258.875 4258.613 4258.613 4258.409 4258.298	9 2 4 3 8 -2 -3 2 -2 0 -1	4259.278 4259.220 4259.163 4259.007 4259.037 4258.896 4258.820 4258.728 4258.638 4258.535 4258.434	-2 -10 -9 -11 -7 1 7 3 6 0 2	8.5 9.5 10.5 11.5 12.5 13.5 14.5 15.5 16.5 17.5 18.5 19.5 20.5

TABLE 1—Continued

$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$
22.5 do 4263.662
74.5 4244.919 -5 4245.402 4 74.5 75.5 4244.534 -7 4245.025 4 75.5 76.5 4244.148 -6 4244.647 7 76.5
75.5 4244.534 -7 4245.025 4 75.5

TABLE 1—Continued

J	R	О-С	P	О-С	Q	О-С	R	О-С	P	О-С	Q	O-C	J
		$^4\Delta_{3/}$	₂ - ⁴ Φ _{5/2}	0-0					$^{4}\Delta_{3/2}$ -	$^4\Phi_{5/2}$	1-1		
8.5 9.5 10.5 11.5 12.5 13.5 14.5 16.5 17.5 18.5 19.5 20.5 22.5 23.5 24.5 24.5 28.5 30.5 31.5 33.5 34.5 35.5 40.5 41.5 42.5 54.5 55.5 56.5 57.5 58.5 57.5 58.5 57.5 58.5 57.5 58.5 57	4205.931 4206.092 4206.251 4206.406 4206.548 4206.687 4207.187 4207.299 4207.403 4207.403 4207.797 4207.796 4207.858 4207.975 4208.005 4207.975 4208.005 4207.975 4207.976 4207.976 4207.976 4207.976 4207.976 4207.975 4207.975 4207.974 4207.975 4207.975 4207.974 4207.975 4207.975 4207.974 4207.975 4207.975 4207.975 4207.975 4207.974	-8 -9 -6 -2 -4 -1 -1 3 3 7 9 4 0 4 -3 0 4 -1 2 -2 5 -1 1 4 4 4 -1 4 -1 4 -1 4 -1 1 2 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	4201.201 4200.900 4200.609 4200.304 4199.694 4199.368 4199.368 4198.540 4197.702 4198.540 4197.344 4196.985 4195.081 4194.678 4195.081 4194.678 4192.984 4192.984 4192.984 4192.984 4192.114 4188.573 4194.678 4194.678 4195.611 4194.678 4195.611 4194.678 4195.611 4196.654 4197.539 4197.539 4197.539 4197.539 4197.539 4197.539 4184.277 4183.572 4184.277 4183.572 4184.277 4183.572 4184.277 4183.572 4184.277 4183.572 4184.277 4183.572 4184.277 4183.572 4184.277 4183.572 4184.279 4184.279 4185.594 4178.704	4 -7 -3 -7 2 -6 -3 0 3 11 8 4 6 8 5 5 4 3 2 0 -1 -1 -1 -4 -5 -4 -4 -6 -5 -4 -2 0 0 2 6 6 6 5 6 0 -2 8 -3	4203.190 4203.138 4203.066 4202.997 4202.930 4202.851 4202.768 4202.270 4202.150 4202.021 4201.890 4201.751 4201.604 4201.450 4201.289 4201.118 4200.565 4200.364 4200.565 4200.364 4200.151 4199.927 4199.694 4199.454 4199.454 4199.454 4199.454 4199.459 4199.459 4199.459 4199.459 4199.459 4199.459 4199.459 4199.3330 4197.330 4191.3330 4192.735 4195.388 4194.918 4193.890 4193.330 4192.735 4195.388 4194.419 4193.890 4193.330 4192.735 4195.388 4194.419 4193.890 4193.330 4192.735 4192.106 4191.438 4190.739 4189.994 4189.213	0 3 -9 -11 -6 -7 -5 -4 4 -1 3 5 5 3 5 4 4 4 2 0 2 1 -1 -3 -6 -7 -5 -5 -6 -7 -4 -4 -3 -1 1 2 4 7 5 6 5 4 2 3 -5 -5 -6 -5 -5 -6	4188.892 4189.046 4189.191 4189.331 4189.464 4189.595 4189.945 4190.044 4190.152 4190.247 4190.355 4190.422 4190.571 4190.654 4190.767 4190.870 4190.914 4190.914 4191.022 4191.049 4191.075 4191.049 4191.075 4191.049 4191.075	1 2 1 0 -2 0 5 -5 -1 -8 0 0 -1 1 2 2 -3 11 1 0 -1 1 2 2 4 12	4183.907 4183.603 4183.603 4183.292 4182.662 4182.334 4182.000 4181.668 4181.312 4180.967 4180.614 4180.249 4179.886 4179.522 4179.135 4177.579 4176.779 4176.362 4175.948 4175.104 4174.245 4173.803 4172.915 4172.465 4172.915 4170.622 4170.150 4170.622 4170.150 4170.622 4170.150 4166.746 4164.713 4165.734 4165.734 4165.734 4165.226 4164.710 4164.193 4163.669 4163.669 4163.669 4163.669 4163.699 4161.596 4161.596 4162.079 4159.905 4159.344 4158.790 4158.228	752 4206-6-30-5-2-5-603-2-4 4-2-1 115121221320-1-2-12-3-1-2-13-5-13-5-13-5	4185, 808 4185,716 4185,611 4185,513 4185,401 4185,159 4185,159 4184,765 4184,620 4184,472 4184,319 4184,160 4183,994 4183,654 4183,474 4183,292 4182,716 4182,515 4182,306 4182,102 4181,887 4181,668 4181,448 4181,221 4180,990 4180,757 4180,519 4174,898 4177,910 4177,628 4179,777 4179,522 4179,777 4179,522 4179,777 4179,522 4179,777 4179,522 4179,777 4179,522 4179,777 4179,522 4179,777 4179,522 4179,777 4179,522 4179,777 4179,522 4179,773 4170,522 4179,773 4170,522 4179,773 4170,522 4179,773 4170,522 4179,773 4170,522 4173,752 4173,578 4174,898 4174,574 4173,216 4173,216 4173,216 4173,216 4173,216 4172,543 4172,543 4172,543 4172,543 4172,543 4172,543 4172,543 4172,543 4171,829	1 1 -5 1 0 1 -5 0 1 -1 -2 -2 0 -1 -3 -3 -1 -2 1 0 0 0 -1 -1 -1 -1 -1 -1 -1 -1 -1 -2 1 2 2 1 3 2 0 0 -1 -2 -3 -11	8.5 9.5 10.5 11.5 12.5 13.5 14.5 16.5 17.5 18.5 19.5 20.5 22.5 23.5 24.5 24.5 24.5 24.5 24.5 33.5 34.5 35.5 36.5 37.5 38.5 37.5 38.5 40.5 41.5 51.5 55.5 56.5 57

TABLE 1—Continued

J	R	О-С	P	О-С	Q	О-С	R	О-С	P	О-С	Q	О-С	J
		$^4\Delta_{3/}$	₂ - ⁴ Φ _{5/2}	2-2					$^4\Delta_{5/2}$ -	$^4\Phi_{7/2}$	0-0		
3.5 4.5 5.5 6.5 7.5 6.5 7.5 10.5 11.5 12.5 13.5 14.5 15.5 16.5 17.5 18.5 19.5 20.5 21.5 22.5 23.5 24.5 25.5 26.5 27.5 28.5 30.5 31.5 3	4164.035 4164.220 4164.376 4164.735 4164.896 4165.055 4165.202 4165.353 4165.764 4166.134 4166.245 4166.361 4166.467 4166.562 4166.562 4166.746 4166.980 4167.058 4167.058 4167.241 4167.240 4167.476 4167.476 4167.476	-3 0 0 0 3 1 0 2 -4 -1 2 2 2 -2 -1 -1 0 0 -1 -1 1 0 0 -2 2 -3 0 -1 1	4160.477 4160.206 4159.947 4159.668 4159.391 4159.112 4158.819 4157.619 4157.619 4157.6391 4156.672 4156.672 4156.672 4156.333 4154.990 4154.641 4154.287 4153.927 4153.564 4153.193 4152.442 4152.057 4151.667 4151.275 4150.877 4150.069 4149.659 4149.659 4149.659 4149.659 4141.927 4141.927 4141.927 4141.927 4141.927 4141.927 4141.928 4140.419 4139.904 4139.386	1 -6 3 -2 -1 3 -2 -7 -2 -9 -2 -1 -1 1 2 0 -1 -1 1 2 0 0 1 0 0 0 -1 -1 1 -1	4161.727 4161.683 4161.646 4161.599 4161.546 4161.495 4161.233 4161.234 4161.234 4161.256 4160.877 4160.776 4160.668 4160.300 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4159.810 4155.819 4155.383 4157.787 4156.961 4156.91 4156.91 4156.517 4155.819 4155.577 4156.91 4156.517 4156.91 4156.517 4155.819 4154.813 4157.789 4157.385 4157.789 4154.813 4157.385 4157.75 4156.961 4156.517 4155.819 4154.555 4155.819 4155.577 4155.819 4154.555 4155.333 4155.076 4154.813 4153.451 4152.872 4152.870 4151.963	4 -3 2 1 -2 3 1 -1 -1 2 -1 0 1 4 2 2 -2 1 1 1 0 2 0 1 1 0 0 -1 1 0 0 -1 1 0 0 -1 1 0 0 -1 1 0 0 -1 1 0 0 -1 1 -3 0 3	4059.884 4060.083 4060.284 4060.6588 4060.6588 4060.195 4061.367 4061.367 4061.534 4062.011 4062.159 4062.304 4062.448 4062.718 4062.972 4063.094 4063.211 4063.323 4063.432 4063.639 4063.639 4063.732 4064.076 4064.076 4064.076 4064.076 4064.576 4064.576 4064.576 4064.601 4064.601 4064.670	2 7 3 2 3 2 1 2 2 2 3 3 3 1 0 1 1 0 0 -1 -1 0 0 -1 -1 1 -2 1 -2 2 -2 -2 -1 -1 -3 -3 1 4 0 0 5 -4 -3 -6	4054.443 4054.153 4053.849 4053.551 4053.244 4052.933 4052.614 4052.301 4051.977 4051.650 4051.318 4050.984 4050.641 4050.300 4049.953 4049.600 4049.245 4048.154 4047.781 4047.403 4047.025 4046.639 4047.403 4047.025 4046.639 4042.569 4041.269 4042.569 4041.269 4041.269 4041.269 4041.269 4040.380 4039.910 4040.380 4039.910 4038.554 4038.662 4040.380 4039.916 4038.554 4038.662 4049.812 4036.662 4040.380 4039.916 4038.554 4038.662 4049.826 4049.826 4049.826 4049.826 4049.826 4049.826 4049.826 4049.826 4049.826 4049.826 4049.826 4031.121 4030.591 4031.647	-1 3 -3 2 1 1 -4 2 1 1 -1 1 -3 0 0 -2 -1 -2 -1 0 -1 -2 -1 -2 -1 -2 -1 -1 1 0 0 1 1 1 1 2 1 0 0 2 3 1 1 1 2 -1 -3 2 3 1	4057,988 4057,965 4057,941 4057,915 4057,878 4057,878 4057,876 4057,670 4057,670 4057,670 4057,492 4057,352 4057,278 4057,195 4057,195 4057,191 4056,933 4056,632 4056,632 4056,518 4056,406 4056,287 4055,578 4055,912 4055,912 4055,578 4055,406 4056,406 4054,888 4054,211 4055,351 4051,368 4054,386 4054,386 4054,386 4054,386 4054,386 4054,386 4054,388 4054,211 4054,032 4053,849 4053,659 4053,679 4051,588 4054,388 4054,211 4054,032 4053,849 4053,659 4053,679 4051,551 4051,588 4054,694 4052,229 4052,007 4051,581 4051,880 4047,808 4047,808 4047,809 4047,007 4046,703 4046,703 4046,703 4046,394	5 1 -1 0 -6 0 4 1 2 1 0 0 -1 -2 1 2 -1 0 1 2 0 -2 -3 -2 0 -2 -1 -1 -1 -1 -1 1 0 0 1 1 1 1 1 2 2 4 1 2 2 3 3 1 3 1 2 1	3.5 4.5 5.5 6.5 7.5 6.5 7.5 8.5 10.5 11.5 12.5 13.5 14.5 15.5 16.5 17.5 18.5 19.5 20.5 21.5 22.5 23.5 24.5 25.5 26.5 27.5 28.5 29.5 31.5 33.5 40.5 41.5 50

TABLE 1—Continued

J R	О-С	P	О-С	Q	О-С	R	О-С	P	О-С	Q	О-С	J
75.5 76.5 77.5 78.5 79.5 80.5 81.5 82.5 83.5 84.5 85.5 86.5 87.5 88.5 89.5 90.5 91.5 92.5 93.5 94.5 95.5 96.5 97.5						4063.949 4063.865 4063.778 4063.682 4063.581 4063.481 4063.369 4063.152 4062.905 4062.781 4062.905 4062.515 4062.317 4062.234 4061.941 4061.786 4061.786 4061.780 4061.781 4061.784	-1 2 4 2 -1 1 -5 -3 4 -9 -3 -1 -1 -2 -1 1 0 1 3 11 4	4028.442 4027.900 4027.346 4026.789 4026.228 4025.661 4025.093 4024.523 4023.348 4022.787 4022.198 4021.608 4021.013 4020.419 4019.816 4019.211 4018.606 4017.993 4017.993 4017.383 4016.767	0 6 3 2 1 -2 -3 -2 -1 -3 -2 -4 -3 -2 2 1 6 10	4046.079 4045.765 4045.411 4045.118 4044.789 4044.456 4044.118 4043.776 4043.434 4043.082 4042.729 4042.375 4042.016 4041.651	-2 1 -2 0 -1 -1 -2 -3 -1 -3 -4 -2 0 -1	75.5 76.5 77.5 78.5 79.5 80.5 81.5 82.5 83.5 84.5 85.5 86.5 87.5 89.5 90.5 91.5 92.5 93.5 94.5 95.5 96.5
	$^4\Delta_{5/2}$	2 - ⁴ Φ _{7/2}	1-1					$^4\Delta_{5/2}$ -	⁴ Ф _{7/2}	2-2		
4.5 5.5 6.5 7.5 8.5 9.5 10.5 11.5 12.5 13.5 4038.386 14.5 4038.554 15.5 4038.720 16.5 4039.202 19.5 4039.365 20.5 4039.508 21.5 4039.656 22.5 4039.798 23.5 4039.930 24.5 4040.431 28.5 29.5 4040.431 28.5 29.5 4040.431 28.5 29.5 4040.656 30.5 4040.762 31.5 32.5 4040.762 31.5 32.5 4040.865 33.5 4041.146 35.5 4041.235 36.5 4041.235 36.5 4041.312 37.5 4041.387 38.5 4041.462 39.5 4041.597 41.5 4041.597 41.5 4041.651 42.5 4041.706 43.5 4041.761 44.5 4041.809 45.5 4041.893 47.5 4041.980 50.5	-1 -6 -10 -6 -11 -9 2 -4 1 2 -1 0 1 5 1 -1 2 1 3 3 -3 -3 0 1 -4 3 1 1 -1	4034.303 4034.048 4033.796 4033.515 4033.248 4032.708 4032.134 4031.841 4031.557 4031.252 4030.953 4030.329 4030.016 4029.695 4029.376 4029.376 4029.376 4029.376 4027.346 4026.282 4026.282 4026.282 4026.282 4026.282 4026.282 4027.346 4028.517 4028.547 4028.547 4028.547 4021.637 4022.852 4022.452 4022.452 4022.452 4022.452 4022.637 4021.637	1 2 11 -5 -2 -3 8 2 -1 9 3 5 -2 -1 0 -1 3 -1 1 -1 0 2 2 -3 1 2 2 1 -1 3 1 1 1 2 1 -2 -1 0 0 0 0 -1 0 -2 -1	4035.244 4035.199 4035.150 4035.150 4035.150 4035.108 4034.934 4034.934 4034.871 4034.653 4034.463 4034.400 4034.303 4034.207 4034.106 4033.998 4033.886 4033.3770 4034.303 4034.207 4033.274 4033.3405 4033.274 4033.2708 4032.556 4032.708 4032.556 4032.708 4032.556 4031.376 4031.908 4031.557 4031.191 4032.075 4031.191 4032.075 4031.191 4031.001 4030.807 4030.613 4030.408 4030.408 4030.408 4030.408 4030.408 4030.408 4030.408	7 1 -4 2 -7 0 -3 -1 1 2 0 -1 2 -1 0 0 -1 1 2 2 0 0 -1 2 0 0 -1 2 0 0 -1 2 0 0 -1 2 0 -1 -1	4014.072 4014.272 4014.480 4014.672 4014.878 4015.071 4015.252 4015.800 4015.979 4016.145 4016.306 4016.466 4016.618 4016.767 4017.320 4017.464 4017.579 4017.702 4017.823 4017.823 4018.329 4018.329 4018.329 4018.419 4018.574 4018.574 4018.574 4018.574 4018.574 4018.574 4018.655 4018.737 4018.851 4019.008 4019.048	1 -9 -6 -14 -5 -4 -12 -11 -3 -4 4 3 1 2 0 -2 -3 0 -7 8 -2 1 6 2 2 0 3 -10 3 -10 -4 8 -2 -4 1 1 2 3	4011.495 4011.231 4010.968 4010.424 4010.146 4009.867 4009.577 4009.577 4009.004 4008.705 4008.395 4008.094 4007.777 4007.466 4007.145 4006.824 4006.497 4006.163 4005.828 4005.142 4004.795 4004.795 4003.356 4002.989 4002.614 4002.234 4001.849 4001.849 4001.85 4000.667 4000.265 3999.857 3999.448 3999.030 3998.174 3997.297 3996.858	6 3 4 1 1 0 2 2 -3 8 7 -1 4 -2 1 1 1 0 1 -3 1 -2 1 1 1 0 1 1 -2 -2 -1 -1 3 3 -1 0 1 1 -3 4	4012.793 4012.767 4012.743 4012.710 4012.629 4012.580 4012.338 4012.311 4012.236 4012.313 4012.083 4011.918 4011.631 4011.631 4011.527 4011.419 4011.308 4011.189 4011.073 4010.945 4010.678 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.544 4010.545 4010.545 4010.540 4009.943 4009.943 4009.943 4009.588 4009.088 4008.901 4008.705 4008.515 4008.308 4007.882 4007.660 4007.437 4007.203	2 -1 -1 -4 -8 -2 2 -1 -1 7 1 -2 0 2 -1 1 1 0 1 -1 -1 2 0 0 -4 1 0 1 -1 -1 0 0 1 1 -1 -1 0 0 0 1 0 0 1 0 0 0 0	4.5 5.5 6.5 7.5 8.5 9.5 10.5 11.5 12.5 13.5 14.5 16.5 17.5 20.5 22.5 22.5 23.5 24.5 24.5 24.5 28.5 30.5 31.5 33.5 34.5 33.5 34.5 35.5 40.5 41.5 42.5 43.5 44.5 45.5 46.5 47.5 48.5 47.5 48.5 58.5

TABLE 1—Continued

J R	О-С	P	О-С	0	О-С	R	О-С	P	О-С	0	О-С	
51.5 52.5 53.5 54.5 55.5 56.5 57.5 58.5 59.5 60.5 61.5 62.5 63.5 64.5 65.5 66.5 67.5 68.5 70.5 71.5 72.5	O-C	4017.766 4017.317 4016.861 4016.401 4015.938 4015.471 4015.000 4014.524 4014.043 4013.559 4013.070 4012.580 4012.083 4011.073 4010.568 4010.050 4009.531 4009.004	O-C -2 -1 -1 -2 -2 -1 -1 -1 -1 -1 -2 2 1 -1 5 2 -1 -1	4029.776 4029.556 4029.334 4029.111 4028.877 4028.640 4028.157 4027.655 4027.398 4026.602 4026.328 4026.047 4025.765 4025.476 4025.178 4024.883 4024.883 4024.883 4024.883 4024.883 4024.883 4024.883 4024.883 4024.883	O-C -2 -3 -2 2 0 -2 1 0 -8 0 0 2 3 3 3 2 4 3 -2 0 -1 -4 -14	R	O-C	P	O-C	Q 4006.968	3	51.5 52.5 53.5 55.5 55.5 56.5 57.5 60.5 61.5 63.5 64.5 66.5 67.5 68.5 69.5 70.5 71.5 72.5 73.5
	⁴ Δ _{7/}	₂ - ⁴ Φ _{9/2}	0-0					⁴ Δ _{7/2} - ⁴	$\Phi_{9/2}$	1-1		
8.5 9.5 3892.694 10.5 3893.077 12.5 3893.077 12.5 3893.439 14.5 3893.619 15.5 3893.793 16.5 3894.130 18.5 3894.130 18.5 3894.445 20.5 3894.600 21.5 3894.445 22.5 3894.895 23.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 25.5 3895.176 26.5 3895.176 27.5 3896.173 37.5 3896.687 29.5 3896.133 33.5 3896.238 34.5 3896.238 34.5 3896.238 34.5 3896.338 34.5 3896.338 34.5 3896.694 38.5 3896.604 38.5 3896.604 38.5 3896.604 38.5 3896.604 38.5 3896.604 38.5 3896.604 38.5 3896.764 40.5 3897.194 45.5 3897.194 45.5 3897.194 45.5 3897.194 45.5 3897.194 45.5 3897.295 48.5 3897.285 49.5 3897.390 52.5 3897.390 52.5 3897.343	0 -6 1 1 -3 0 1 0 4 1 0 0 1 1 1 3 3 2 2 3 0 2 2 2 1 -1 0 0 -3 -3 -4 2 3 3 -2 -3 0	3888.262 3887.995 3887.413 3887.433 3887.153 3886.562 3886.562 3886.265 3885.966 3885.352 3885.042 3884.725 3884.082 3883.751 3883.422 3882.747 3882.404 3882.056 3881.706 3881.706 3881.350 3880.263 3879.513 3879.513 3879.513 3879.513 3879.513 3879.513 3879.513 3879.513 3879.513 3879.513 3877.977 3877.582 3877.185 3877.977 3877.582 3877.185 3874.289 3873.861 3873.427 3873.861 3873.427 3872.550 3872.105 3871.658	-1 5 -3 2 7 -5 -2 -3 0	3890.257 3890.223 3890.178 3890.133 3890.083 3890.027 3889.971 3889.971 3889.546 3889.703 3889.546 3889.546 3889.371 3889.277 3889.371 3889.277 3889.813 3888.65 3888.754 3888.65 3888.754 3888.65 3888.754 3888.754 3888.754 3888.754 3888.65 3888.754 3888.65 3888.754 3888.65 3888.754 3888.65 3888.754 3888.65 3888.754 3888.65 3888.756 3887.710 3887.563 3887.563 3887.563 3887.563 3887.563 3887.563 3887.563 3887.563 3887.563 3887.563	-6 -1 -4 -3 -3 -5 -3 -3 -4 -1 0 0 2 2 3 2 1 2 2 2 1 2 2 2 1 1 0 0 1 1 1 0 0 1 1 1 0 0 1 1 1 1	3870.517 3870.686 3870.852 3871.012 3871.181 3871.330 3871.481 3871.909 3872.046 3872.172 3872.298 3872.426 3872.550 3872.666 3872.780 3872.884 3872.992	1 -1 -1 -4 7 1 3 -3 1 2 4 4 -1 -3 1 5 5 5 -1 0	3864.734 3864.458 3864.458 3863.898 3863.620 3863.319 3863.028 3862.423 3862.423 3862.423 3862.423 3860.460 3860.533 3860.208 3859.873 3859.873 3859.873 3859.540 3858.855 3857.452 385	-1035	3866.958 3866.914 3866.824 3866.772 3866.716 3866.653 3866.583 3866.521 3866.448 3866.211 3866.211 3866.211 3866.211 3865.940 3865.842 3865.520 3865.842 3865.520 3865.412 3865.520 3865.412 3864.927 3864.927 3864.927 3864.392 3863.363 3863.662 3863.363 3863.662 3863.373 3863.662 3863.373 3862.2731 3862.578 3862.578 3862.578 3862.578 3862.623 3863.813	-1038	

TABLE 1—Continued

J	R	О-С	P	О-С	Q	О-С	R	О-С	P	О-С	Q	О-С	J
55.5			3871.206	-1	3884.222	-1			3847.566*	-1066			55.5
56.5			3870.750	o	3884.000	i			3847.142*		3860.332*	-1085	56.5
57.5			3870.291	ŏ	3883.775	$\hat{2}$			3846.721*		3860.139*		57.5
58.5			3869.828	Ŏ	3883.544	1			3846.292*		3859.943*		58.5
59.5			3869.363	2	3883.312	3			3845.862*	-1136	3859.747*		59.5
60.5			3868.894	3	3883.077	6			3845.428*	-1153	3859.540*		60.5
61.5			3868.420	3	3882.834	4			3845.001*	-1157	3859.348*	-1159	61.5
62.5			3867.944	4	3882.590	4			3844.564*	-1165	3859.144*	-1166	62.5
63.5			3867.464	4	3882.342	3			3844.126*	-1167	3858.942*	-1162	63.5
64.5			3866.982	5	3882.095	7					3858.740*	-1151	64.5
65.5			3866.496	5	3881.839	4					3858.536*	-1132	65.5
66.5			3866.005	3	3881.580	2					3858.333*	-1101	66.5
67.5			3865.520	8	3881.325	4					3858.129*	-1057	67.5
68.5			3865.020	1	3881.062	2					3857.927*	-997	68.5
69.5			3864.529	4	3880.800	3					3857.718*	-926	69.5
70.5			3864.028	-1	3880.531	-2					3857.516*	-827	70.5
71.5			3863.532	0	3880.263	-4					3857.306*	-716	71.5
72.5			3863.028	-5	3879.998	-3					3857.093*	-581	72.5
73.5			3862.529	-6	3879.729	-5					3856.869*	-429	73.5
74.5			3862.031	-5	3879.461	-5					3856.665*	-224	74.5
75.5			3861.528	-9	3879.191	-9					3856.462*	18	75.5
76.5			3861.031	-9	3878.929	-5					3856.260*	301	76.5
77.5			3860.533	-12	3878.664	-6					3856.038*	609	77.5
78.5					3878.400	-8					3855.830*	981	78.5
79.5					3878.148	0					3855.624*	1411	79.5
80.5					3877.898	6					3855.425*	1909	80.5
81.5					3877.647	7							81.5
82.5					3877.404	10							82.5

 $^4\Delta_{7/2}$ $^{-4}\Phi_{9/2}$ 2-2

J	R	О-С	P	О-С	Q	О-С	J	R	О-С	P	О-С	Q	О-С
19.5					3842.981	13	56.5			3823.822	1	3836.949	1
20.5					3842.888	9	57.5			3823.351	2	3836.710	3
21.5	3848.032	5			3842.796	10	58.5			3822.875	1	3836.466	4
22.5	3848.168	5	3837.448	-1	3842.694	4	59.5			3822.398	5	3836.219	6
23.5	3848.298	4	3837.120	4	3842.594	6	60.5			3821.913	3	3835.964	4
24.5	3848.424	2	3836.780	2	3842.485	2	61.5			3821.426	5	3835.709	7
25.5	3848.548	3	3836.440	4	3842.371	-3	62.5			3820.937	8	3835.449	8
26.5	3848.662	-2	3836.087	-2	3842.259	-1	63.5			3820.442	9	3835.193*	19
27.5	3848.779	0	3835.735	-4	3842.142	-1	64.5			3819.951*	19		
28.5	3848.888	-1	3835.382	-3	3842.015	-5	65.5					3834.606*	-23
29.5	3848.994	-2			3841.892	-2	66.5					3834.334*	-17
30.5	3849.098	0	3834.660	-3	3841.762	-2	67.5					3834.059	-8
31.5			3834.293	-4	3841.627	-3	68.5					3833.772	-7
32.5			3833.923	-3	3841.487	-5	69.5					3833.480	-7
33.5	3849.378	-2	3833.547	-4	3841.348	-1	70.5					3833.184	-7
34.5	3849.460	-5			3841.201	-1	71.5					3832.884	-6
35.5	3849.542	-5	3832.790	1			72.5					3832.578	-7
36.5	3849.618	-6	3832.399	-3	3840.896	-2	73.5					3832.270	-5
37.5	3849.698	0	3832.008	-4	3840.736	-3	74.5			3814.687	-2	3831.957	-4
38.5	3849.767	1	3831.620	3	3840.572	-4	75.5			3814.140	-2		
39.5	3849.832	0	3831.216	-2	3840.408	-1	76.5			3813.591	2	3831.317	-2
40.5	3849.895	2	3830.814	-1	3840.238	0	77.5					3830.989	-3
41.5	3849.950	0	3830.409	1	3840.064	1	78.5			3812.470	-2	3830.661	1
42.5	3850.003	1	3829.997	1	3839.886	2	79.5			3811.904	-2	3830.323	0
43.5	3850.057	6	3829.584	2	3839.700	0	80.5			3811.339	2	3830.005	22
44.5	3850.100	4	3829.164	1	3839.514	1	81.5			3810.765	2	3829.640	2
45.5			3828.739	0	3839.323	1	82.5					3829.291	2
46.5			3828.312	-1	3839.129	3	83.5			3809.596	-6	3828.940	2 5 3
47.5			3827.881	0	3838.926	-1	84.5			3809.022	5	3828.580	3
48.5			3827.448	2	3838.726	3	85.5			3808.427	1	3828.224	8
49.5			3827.008	0	3838.518	2	86.5			3807.834	2	3827.854	5
50.5			3826.563	-1	3838.308	4	87.5				_	3827.476	-3
51.5			3826.117	-1			88.5			3806.634	2	3827.108	2
52.5			3825.665	-1		_	89.5			3806.026	1		
53.5			3825.211	0	3837.646	2	90.5			3805.413	-3		
54.5			3824.758	7	3837.419	2	91.5			3804.794	-9		
55.5			3824.290	2	3837.187	3							

TABLE 2

Consequence (in part 1) for the C4 × V4% David of Z-C1

			Rotatio	nnal Constants	Rotational Constants (in cm ⁻¹) for the $C^4\triangle - X^4\Phi$ Bands of ZrCl	e $\mathbf{C}^4 \triangle - \mathbf{X}^4 \Phi \mathbf{B}$	ands of ZrCl			
Transition	ion	T_{w}	$\mathbf{B}_{\mathbf{v}}'$	$10^8 \times D_{\rm v}'$	$10^2 imes ext{p}_{ ext{v}}'$	$10^7 \times \rm p_{Dv}{'}$	$\mathbf{B}_{\mathrm{v}}^{"}$	$10^8 \times D_v^{"}$	$10^{11} \times \mathrm{H_v}''$	$10^{15} \times L_{\rm v}^{"}$
$^4\Delta_{\scriptscriptstyle 1/2}$ $^4\Phi_{\scriptscriptstyle 3/2}$	0-0	4300.38412(50) 0.1171518(47)	0.1171518(47)	5.376(39)	-1.2089(15)	2.671(26)	0.1201246(48) 15.890(63)	15.890(63)	0.9732(88)	-0.3613(51)
	1-1	4280.20175(44) 0.1166470(56)	0.1166470(56)	5.502(54)	-0.8018(16)	0.886(26)	0.1193543(55)	8.327(54)	0.1311(16)	!
	2-2	4259.50793(88) 0.1161193(93)	0.1161193(93)	5.73(23)	-0.63227(93)	ŀ	0.1187165(93)	6.99(24)	0.114(40)	-0.067(33)
$^4\Delta_{^3\mathcal{I}^2}{}^4\Phi_{^{5\mathcal{I}_2}}$	0-0	4203.4231(15)	0.117249(19)	6.06(57)	;	ı	0.120140(19)	9.90(60)	12.936(63)	ţ
	1-1	4186.44147(90)	0.1166852(84)	5.23(17)	ŀ	1	0.1199711(91)	29.01(30)	3.381(65)	-2.015(64)
	2-2	4161.84049(43)	0.1161620(69)	5.37(18)	ı	ı	0.1185871(69)	6.58(18)	0.218(14)	I
$^4\Delta_{5/2}$ $^4\Phi_{7/2}$	0-0	4058.01515(37) 0.1174188(40)	0.1174188(40)	5.354(31)	ŀ	ł	0.1194671(39)	4.792(32)	-0.04289(98)	ł
	1-1	4035.40606(55)	0.1168957(74)	5.53(14)	;	ŀ	0.1189822(74)	5.95(14)	0.0882(74)	;
	2-2	4012.84360(62)	0.116336(12)	4.80(40)	:	;	0.118448(12)	6.27(45)	1.387(59)	!
$^4\Delta_{7/2}$ $^4\Phi_{9/2}$	0-0	3890.42399(73)	0.1175978(77)	5.50(13)	1	1	0.1196049(81)	8.43(19)	1.013(31)	-1.206(24)
	1-1	3867.1431(12)	0.1170502(72)	5.19(16)	;	ł	0.1189395(89)	-4.66(59)	-7.34(24)	9.57(33)
	2-2	3843.8346(19)	0.1165100(82)	5.349(92)	ł	ł	0.1186911(87)	8.52(17)	0.470(28)	-0.226(16)

Note: The numbers in parentheses are one standard deviation in the last two digits.

side of the 0-0 band of the $^4\Delta_{3/2}-^4\Phi_{5/2}$ subband and these lines also probably belong to another transition. The $^4\Delta_{3/2}-^4\Phi_{5/2}$ subband was difficult to assign because of a global perturbation of the $^4\Phi_{5/2}$ spin component and the extra lines. No extra heads have been observed in the regions of the $^4\Delta_{5/2}-^4\Phi_{7/2}$ and $^4\Delta_{7/2}-^4\Phi_{9/2}$ subbands. The 0-0, 1-1, and 2-2 bands of all the subbands were rotationally analyzed. A part of the structure of the 0-0 band of the $^4\Delta_{7/2}-^4\Phi_{9/2}$ subband is presented in Fig. 2.

As is often the case, no transitions having $\Delta\Sigma \neq 0$ were observed and the spin-orbit intervals could not be determined directly. The subbands of the different spin components were fitted separately using a simple effective term energy expression:

$$F_{v}(J) = T_{v} + B_{v}J(J+1) - D_{v}[J(J+1)]^{2}$$

$$+ H_{v}[J(J+1)]^{3} + L_{v}[J(J+1)]^{4}$$

$$\pm 1/2[p_{v}(J+1/2) + p_{Dv}(J+1/2)^{3}].$$
[1]

The rotational lines were weighted according to resolution, extent of blending, and the effect of perturbations. Perturbed lines were not included in the fit and the badly blended lines were heavily deweighted. Occasionally higherorder effective rotational constants H_v and L_v are required in some spin components to obtain a satisfactory fit. The higher-order rotational constants clearly have no mechanical meaning but they reflect interactions with the other spin components and other electronic states. The observed line positions for the different subbands are provided in Table 1 and the molecular constants for the different bands are provided in Table 2. The e/f parity assignment in the $^4\Delta_{1/2}$ - $^4\Phi_{3/2}$ subband was made arbitrarily to provide a negative Ω -doubling constant p_v in the $^4\Delta_{1/2}$ state. The higher-order Ω -doubling constant p_{Dv} was also determined for the v=0and 1 vibrational levels of this state.

DISCUSSION

Although the bands of ZrCl in the visible and near-ultraviolet regions have been known for decades, the identity of the ground state is still an open question. The previously suggested $^4\Sigma^-$ ground state is not consistent with the recently observed $X^4\Phi$ ground state of TiF (8) and TiCl (9). We suspect that Jordan *et al.* (2) have, in fact, measured a $^4\Delta-X^4\Phi$ or a $^4\Gamma-X^4\Phi$ transition in the region 400–415 nm. The present $^4\Delta-^4\Phi$ transition of ZrCl is analogous to the $C^4\Delta-X^4\Phi$ transition of TiCl (9). Since no theoretical calculations are available for ZrCl, we have used the recent calculations for the isovalent TiF (6) molecule as a guide. As predicted by Harrison for TiF (6), there are quartet and doublet manifolds of electronic states, of which a $^4\Phi$ state is the lowest in energy. Among the doublet states, a $^2\Phi$ state

is lowest at about 5200 cm⁻¹ above the ground $X^4\Phi$ state. Although the $^2\Phi$ - $X^4\Phi$ separation will be different for TiF and TiCl, the general pattern of energy levels should be similar. For ZrCl, the doublet–quartet separation may well change considerably and it is possible that a $^2\Phi$ or even a $^2\Delta$ state could become the ground state. Further experimental and theoretical work is necessary to confirm the ground state assignment for ZrCl.

Several perturbations have been observed in the spectrum of the ZrCl $C^4\Delta-X^4\Phi$ transition. The excited $C^4\Delta$ state is unperturbed. The ground $X^4\Phi_{3/2}$ spin component with v=1 is affected by a local perturbation at about J''=67.5 and the v=0 level is affected by a global perturbation which shifts $B_{\rm eff}$ to a higher value. The v=0 and 1 vibrational levels of the $X^4\Phi_{5/2}$ spin component are affected by strong global perturbations. Local perturbations are also present in the v=1 vibrational level of the $X^4\Phi_{9/2}$ spin component at J''=52.5 and probably in the v=2 vibrational level at J''=66.5. Clearly there are one or more low-lying electronic states which interact with the $X^4\Phi_{3/2}$, $X^4\Phi_{5/2}$, and $X^4\Phi_{9/2}$ spin components.

In the absence of $\Delta v \neq 0$ bands, only limited vibrational data could be extracted from the present analysis. As seen in Table 2, the excited $^4\Delta_{1/2}$ spin component has large Ω -doubling, presumably because of interaction with another state such as a $^2\Pi$ state. This observation is in contrast to the results for TiCl where no Ω -doubling was observed in the $C^4\Delta$ state. The observation in ZrCl of the transition analogous to the $G^4\Phi$ - $X^4\Phi$ transition of TiF (8) and TiCl (9) will provide more information about the electronic structure. So far this transition has not been detected.

The Hund's case (a) equilibrium constants were derived by averaging the effective rotational constants of each spin component. The relationship between the effective constants and Hund's case (a) constants for $^4\Delta$ states (14) and $^4\Phi$ states (8) is well known. For the $C^4\Delta$ state the $B_{\rm eff}$ values for each vibrational level were averaged and then $B_{\rm e}'$ = $0.11769 \text{ cm}^{-1} \text{ and } \alpha_{e}' = 0.00054 \text{ cm}^{-1} \text{ were determined. An}$ estimate for the spin-orbit constant $A'_0 = 62 \text{ cm}^{-1}$ was made from the different $B_{\rm eff}$ values. The $X^4\Phi$ state shows the effects of both global and local perturbations so the procedure described above was used only for the relatively unaffected $X^4\Phi_{9/2}$ and $X^4\Phi_{7/2}$ spin components. The values of $B_e''=$ 0.1196 cm^{-1} , $\alpha_e'' = 0.0004 \text{ cm}^{-1}$, and $A_0'' = 69 \text{ cm}^{-1}$ were obtained for the $X^4\Phi$ state. The B'_e and B''_e values provide $r'_{\rm e} = 2.3852 \text{ Å}$ and $r''_{\rm e} = 2.3661 \text{ Å}$ for the excited and ground states, respectively.

The electron configurations of the observed states are expected to be the same as those for the isovalent TiCl molecule. Following the simple ionic bonding model, the $C^4\Delta$ and $X^4\Phi$ states of ZrCl arise from $\pi^2\delta^1$ and $\sigma^1\pi^1\delta^1$ configurations on the metal atom, respectively.

CONCLUSIONS

We have recorded the emission spectrum of ZrCl in the region 3000–10 000 cm⁻¹ using a Fourier transform spectrometer. Four groups of prominent bands with the 0–0 R heads at 4305.7, 4208.0, 4064.7, and 3897.4 cm⁻¹ have been assigned as the ${}^4\Delta_{1/2} - {}^4\Phi_{3/2}$, ${}^4\Delta_{3/2} - {}^4\Phi_{5/2}$, ${}^4\Delta_{5/2} - {}^4\Phi_{7/2}$, and ${}^4\Delta_{7/2} - {}^4\Phi_{9/2}$ subbands of the $C^4\Delta - X^4\Phi$ transition analogous to the same transition of TiCl in the region 3100–3400 cm⁻¹ (9). Each of the 0–0 bands is followed by weaker 1–1, 2–2, and 3–3 vibrational bands to lower wavenumbers. A rotational analysis of several of these bands has been obtained and the molecular constants have been determined. The lowest ${}^4\Phi$ state most probably is the ground state of this molecule but we do not have any direct evidence to prove this. Further experimental and theoretical work will be necessary to prove our proposed assignment.

ACKNOWLEDGMENTS

We thank J. Wagner, C. Plymate, and M. Dulick of the National Solar Observatory for assistance in obtaining the spectra. The National Solar Observatory is operated by the Association of Universities for Research in Astronomy, Inc., under contract with the National Science Foundation. The research described here was supported by funding from the NASA labora-

tory astrophysics program. Some support was also provided by the Petroleum Research Fund administered by the American Chemical Society and the Natural Sciences and Engineering Research Council of Canada.

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